Entanglement Distribution with Separable States



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KIAS-2017



A and B are separable



T. S. Cubitt, F. Verstraete, W. Du[°]r, and J. I. Cirac, Phys. Rev. Lett. 91, 037902 (2003).

The point is that the messenger remains separable at all stages.



$$\rho_{ABM} = \sum_{i} p_{i} \rho_{A}^{i} \otimes \rho_{B}^{i} \otimes \rho_{M}^{i}$$

But the initial state should be very peculiar.

$$\rho_{ABM} = \frac{1}{6} \sum_{k=0}^{3} |\psi_k\rangle \langle \psi_k| \otimes |\psi_{-k}\rangle \langle \psi_{-k}| \otimes |0\rangle \langle 0| + \frac{1}{6} \sum_{i=0}^{1} |i, i, 1\rangle \langle i, i, 1| \langle i, i| \langle i,$$

$$\left|\psi_{k}\right\rangle = \frac{1}{\sqrt{2}} \left(\left|0\right\rangle + e^{\frac{i\pi}{2}k}\left|1\right\rangle\right)$$

Our Goal:

To understand this result

Generalization to d-level states



Generalization to GHZ states



Let us start from the very beginning:

$$|\psi_k\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \alpha_k |1\rangle)$$

$$|\phi_k\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \beta_k |1\rangle)$$



After Alice CNOT

$$\left|\Psi_{k}\right\rangle_{AM} = \frac{1}{\sqrt{2}} \left(\left|00\right\rangle + \alpha_{k}\right|11\right)$$



$$|\Psi_k^1\rangle_{AMB} = \frac{1}{2} \left(|00\rangle + \alpha_k|11\right) \left(|0\rangle + \beta_k|1\rangle\right)$$

The messenger is sent to Bob:

 $|\Psi_k^1\rangle_{AMB} = \frac{1}{2} \left(|000\rangle + \alpha_k|110\rangle + \beta_k|001\rangle + \alpha_k\beta_k|111\rangle\right)$

We rearrange the indices for simplicity

 $|\Psi_k^1\rangle_{ABM} = \frac{1}{2} \left(|000\rangle + \alpha_k|101\rangle + \beta_k|010\rangle + \alpha_k\beta_k|111\rangle\right)$



From the previous page:

 $|\Psi_k^1\rangle_{ABM} = \frac{1}{2} \left(|000\rangle + \alpha_k|101\rangle + \beta_k|010\rangle + \alpha_k\beta_k|111\rangle\right)$

After Bob CNOT operation, the state is:

 $|\Psi_k^2\rangle_{ABM} = \frac{1}{2} \left(|000\rangle + \alpha_k|101\rangle + \beta_k|011\rangle + |110\rangle\right)$

$$=\frac{1}{\sqrt{2}}\left(|\phi_{+}\rangle|0\rangle_{M}+|\chi_{k}\rangle|1\rangle_{M}\right)$$

$$\chi_k \rangle = \frac{1}{\sqrt{2}} (\alpha_k |10\rangle + \beta_k |01\rangle)$$

However the particle M has entangled itself

with A and B

in both stages of the process.

In stage 1: After Alice operation:



$$|\Psi_k^1\rangle_{AMB} = \frac{1}{2} \left(|00\rangle + \alpha_k|11\right) \left(|0\rangle + \beta_k|1\rangle\right)$$

In stage 2: After Bob's Operation



$$|\Psi_k^2\rangle_{AMB} = \frac{1}{\sqrt{2}} \left(|\phi_+\rangle|0\rangle_M + |\chi_k\rangle|1\rangle_M \right)$$

How to remove the entanglement?

For (1) we use Symmetrization For (2) we use Mixing



 $|\phi_{+}\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) \qquad \qquad |\phi_{-}\rangle = \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$

$$\frac{1}{2} \left(\left| \phi_{+} \right\rangle \left\langle \phi_{+} \right| + \left| \phi_{-} \right\rangle \left\langle \phi_{-} \right| \right) = \frac{1}{2} \left(\left| 00 \right\rangle \left\langle 00 \right| + \left| 11 \right\rangle \left\langle 11 \right| \right) \right)$$



$$|\Psi_k^2\rangle_{ABM} = \frac{1}{\sqrt{2}} \left(|\phi_+\rangle|0\rangle + |\chi_k\rangle|1\rangle \right)$$

$$\rho_{_{ABM}}^2 = \sum_k |\Psi_k^2\rangle \langle \Psi_k^2|$$

$$\rho_{ABM}^{2} = \sum_{k} \frac{1}{2} \left(|\phi_{+}\rangle \langle \phi_{+}| \otimes |0\rangle \langle 0| + |\chi_{k}\rangle \langle \chi_{k}| \otimes |1\rangle \langle 1| \right) \\ + \frac{1}{2} \left(|\phi_{+}\rangle \langle \chi_{k}| \otimes |0\rangle \langle 1| + |\chi_{k}\rangle \langle \phi_{+}| \otimes |1\rangle \langle 0| \right)$$

We demand that;

$$\sum_{k} |\chi_{k}\rangle = 0$$

$$|\chi_k\rangle = \frac{1}{\sqrt{2}} (\alpha_k |10\rangle + \beta_k |01\rangle)$$

$$\sum_{k} \alpha_{k} = 0$$

So after Bob's operation we have this:







$$\rho_{ABM}^2 = \sum_k \frac{1}{2} \left(|\phi_+\rangle \langle \phi_+| \otimes |0\rangle \langle 0| + |\chi_k\rangle \langle \chi_k| \otimes |1\rangle \langle 1| \right)$$

(1): Removing entanglement in the first stage



$$\begin{split} |\Psi_k^1\rangle_{AMB} &= \frac{1}{2} \left(|00\rangle + \alpha_k |11\rangle \right) \left(|0\rangle + \beta_k |1\rangle \right) \\ |\Psi_k^1\rangle_{AMB} &= \frac{1}{2} \left(|000\rangle + \alpha_k |110\rangle + \beta_k |001\rangle + |111\rangle \right) \\ |\Psi_k^1\rangle_{AMB} &= \frac{1}{2} \left(\sqrt{2} |GHZ\rangle + \alpha_k |110\rangle + \beta_k |001\rangle \right) \end{split}$$



B is obviously separate from A and M.

So If we make the state symmetric with respect to B and M,

then M will be separate from A and B.

$|\Psi_k^1\rangle_{AMB} = \frac{1}{2} \left(\sqrt{2} |GHZ\rangle + \alpha_k |110\rangle + \beta_k |001\rangle \right)$

We want this to be symmetric:

$$\rho_{_{AMB}} = \sum_{k} |\Psi_{k}^{1}\rangle \langle \Psi_{k}^{1}|$$

The Term |GHZ><110| vanishes due to : The Term |GHZ><001| vanishes due to :



The Term |110><001| vanishes if we assume that:

 $\sum_{k} \alpha_{k}^{2} = 0$

$|\Psi_k^1\rangle_{AMB} = \frac{1}{2} \left(\sqrt{2} |GHZ\rangle + \alpha_k |110\rangle + \beta_k |001\rangle \right)$

But two terms remain which are not symmetric::

 $\pi_{110} + \pi_{001}$

So we add the terms

 $\pi_{101} + \pi_{010}$

This addition does not affect the process?

The Solutions



$$\rho_{ABM} = \frac{1}{6} \sum_{k=0}^{3} |\psi_k\rangle \langle \psi_k| \otimes |\psi_{-k}\rangle \langle \psi_{-k}| \otimes |0\rangle \langle 0| + \frac{1}{6} \sum_{i=0}^{1} |i, i, 1\rangle \langle i, i, 1|$$

$$|\psi_k\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle + e^{\frac{i\pi}{2}k}|1\rangle\right)$$

T. S. Cubitt, F. Verstraete, W. Du[¨]r, and J. I. Cirac, Phys. Rev. Lett. 91, 037902 (2003).



 $CNOT|i,j\rangle = |i,i+j\rangle$

$$|\psi\rangle_{k} = \frac{1}{\sqrt{d}} \left(|0\rangle + \alpha_{k}|1\rangle + \beta_{k}|2\rangle + \cdots \xi_{k}|d-1\rangle\right)$$

How many states we need?

What are the form of these states?

$$|\psi_k\rangle = \frac{1}{\sqrt{d}} \sum_{j=0}^{d-1} \omega^{ks_j} |j\rangle$$

$$\omega = e^{\frac{2\pi i}{K}}$$

$$|\psi_k\rangle = \frac{1}{\sqrt{d}} \left(\omega^{s_0 k} |0\rangle + \omega^{s_1 k} |1\rangle + \omega^{s_2 k} |2\rangle + \cdots \right)$$

The number of states = K





 $s_i + s_j \neq s_k + s_l$

For d=3 (Qutrits)



 $s_i + s_j \neq s_k + s_l$



K=7 is the 7th root of unity.



 $s_i + s_j \neq s_k + s_l$

Distribution of GHZ states

 $\alpha_k \beta_k \gamma_k = 1$



$|\Psi_k^1\rangle_{AM_1M_2BC} = (|000\rangle + \alpha_k|111\rangle) (|0\rangle + \beta_k|1\rangle) (|0\rangle + \gamma_k|1\rangle)$







$|\Psi_k^2\rangle = \frac{1}{2} \left(|00\rangle|GHZ\rangle + |01\rangle|\chi_k^{01}\rangle + |10\rangle|\chi_k^{10}\rangle + |11\rangle|\chi_k^{11}\rangle\right)$





$$|\chi_k^{01}\rangle = \frac{1}{\sqrt{2}} \left(\gamma_k |000\rangle + \gamma_k^{-1} |111\rangle \right)$$
$$|\chi_k^{10}\rangle = \frac{1}{\sqrt{2}} \left(\beta_k |000\rangle + \beta_k^{-1} |111\rangle \right)$$

$$|\chi_k^{11}\rangle = \frac{1}{\sqrt{2}} \left(\alpha_k |000\rangle + \alpha_k^{-1} |111\rangle \right)$$

The results of measurement of the messengers determine the type of GHZ state obtained!

Removing the entanglement after stage 2 by mixing:

$$\sum_{k} \alpha_{k} = 0$$
$$\sum_{k} \beta_{k} = 0$$

 $\sum_k \gamma_k = 0$

$$\sum_{k} \frac{\alpha_{k}}{\beta_{k}} = \sum_{k} \frac{\alpha_{k}}{\gamma_{k}} = \sum_{k} \frac{\beta_{k}}{\gamma_{k}} = 0$$

Symmetrization



All local operations are done by A on the messengers

Messengers are separate from B and C





We make this state symmetric with respect to $(M_1^{\bullet \bullet} B)$ and $(M_2^{\bullet \bullet} C)$

Removing the entanglement after stage 1 by symmetrization

$\sum_{k} \alpha_{k}^{2} = \sum_{k} \beta_{k}^{2} = \sum_{k} \gamma_{k}^{2} = 0$

Removing the entanglement after stage 1 by symmetrization

 $\sum_{k} \alpha_{k} = \sum_{k} \beta_{k} = \sum_{k} \gamma_{k} = 0$ $\sum_{k} \frac{\alpha_{k}}{\beta_{k}} = \sum_{k} \frac{\alpha_{k}}{\gamma_{k}} = \sum_{k} \frac{\beta_{k}}{\gamma_{k}} = 0$

$$\sum_{k} \alpha_{k}^{2} = \sum_{k} \beta_{k}^{2} = \sum_{k} \gamma_{k}^{2} = 0$$

Solution

 $\sum_{k} \alpha_{k} = 0$

 $\sum_{k}\beta_{k}=0 \quad \blacklozenge$



 $\sum_{k} \alpha_{k}^{2} = 0$

 $\sum_{k} \mathbf{p}_{k}^{\underline{A}} = \mathbf{0} \quad \blacklozenge$

 $\sum_{k} q x_{k}^{\mathfrak{D}} = \emptyset \quad \blacklozenge$

 $\sum_{k} \frac{\alpha_{k}}{\beta_{k}} = 0 \quad \blacklozenge$

 $\sum_{k} \frac{\alpha_{k}}{\gamma_{k}} = 0 \quad \blacklozenge$

 $\sum_{k} \frac{\beta_{k}}{\gamma_{k}} = 0 \quad \blacklozenge$

 $\beta_k = \alpha_k^2$ $\gamma_k = \alpha_k^{-3}$

 $\alpha_k \beta_k \gamma_k = 1$

Solution







 $\sum_{k} \alpha_{k}^{6} = 0$



$\alpha_k = \omega^k$

 $1 + \omega^{1} + \omega^{2} + \omega^{3} + \omega^{4} + \omega^{5} + \omega^{6} = 0$ $1 + \omega^{2} + \omega^{4} + \omega^{6} + \omega^{1} + \omega^{3} + \omega^{5} = 0$ $1 + \omega^{3} + \omega^{6} + \omega^{2} + \omega^{5} + \omega^{1} + \omega^{4} = 0$

$\bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet$

$$\omega = e^{\frac{2\pi i}{7}}$$

We can distribute GHZ states with probability

Summary:

There is a systematic way to produce Bell states and GHZ GHZ(n) states in any dimension, by using only separable messengers.

Thanks for your attention