

Heisenberg's Error Disturbance Relation (EDR)

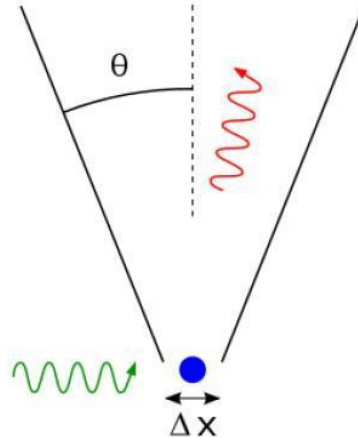
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(Oct. 29, 2013)

Outline

- Heisenberg-Bohr's Microscope
- The difference between preparation variances and EDR.
- Werner's approach to prove the EDR.
- Ozawa's error disturbance relation and disproving the EDR.
- Experimental test of EDR.

Heisenberg-Bohr's microscope



- The uncertainty in position is the smallest distance at which two bodies might be distinguished.

$$\epsilon(x) = \frac{\lambda}{2 \sin \theta}$$

- The disturbance in momentum of the electron is at list equal to the uncertainty in x component of photons momentum.

$$\eta(p_x) = p \sin \theta$$

What is EDR exactly?

- It's impossible to measure the position of a mass without disturbing its momentum,


$\epsilon(A)$ Is the error in measurement, and not the variance of distribution.

$$\epsilon(A)\eta(B) \geq \frac{|\langle [A, B] \rangle|}{2} = \frac{\hbar}{2}.$$

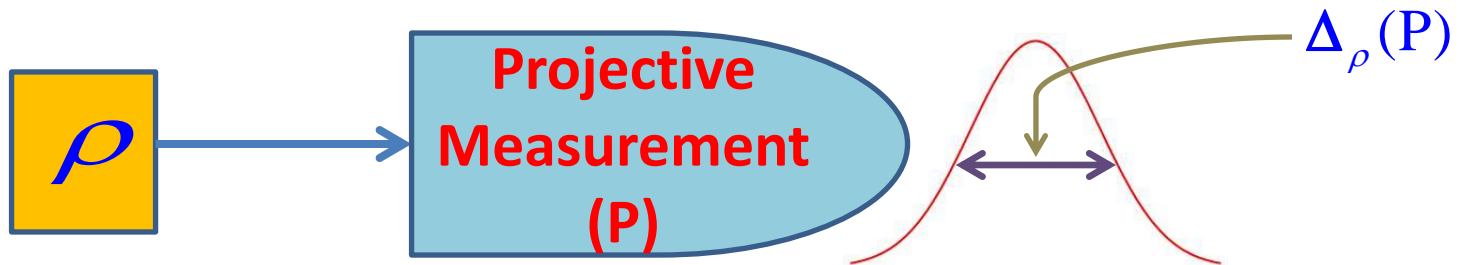
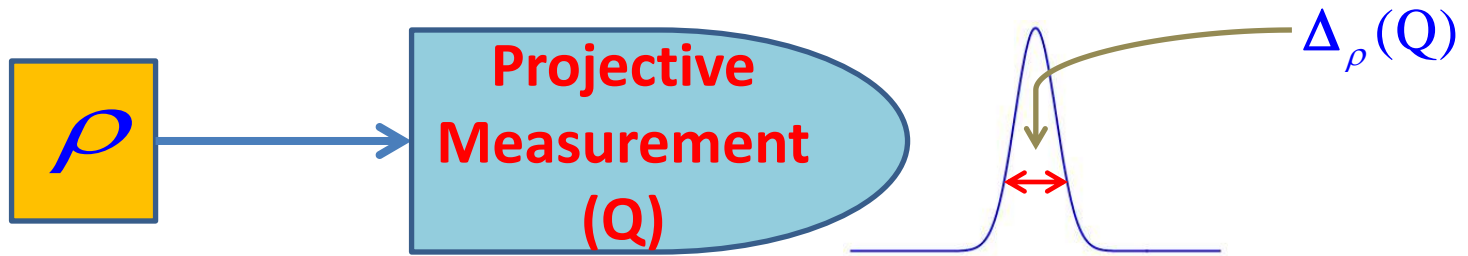
- It's different from the inequality relation between the variances of A and B,

$$\sigma(A)\sigma(B) \geq \frac{|\langle [A, B] \rangle|}{2} = \frac{\hbar}{2}.$$

This inequality does not relate to the measurements.

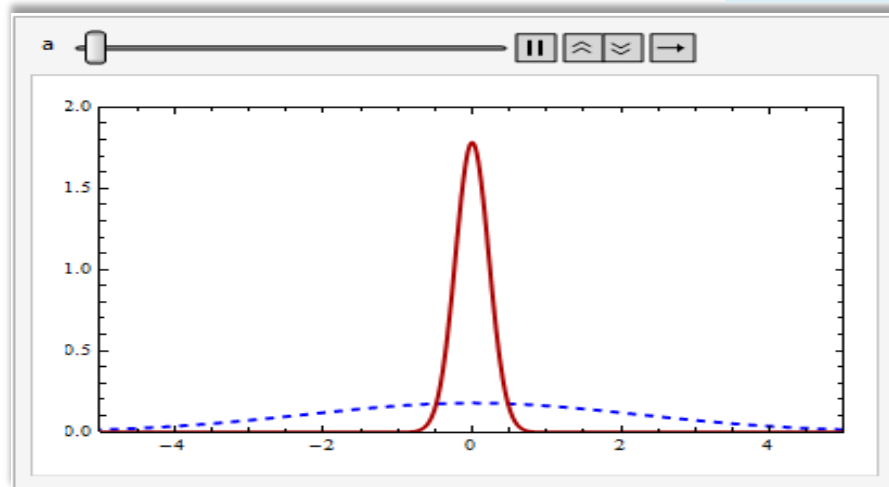

$$\sigma(A) = \langle (A - \langle A \rangle)^2 \rangle^{1/2}$$

- It's impossible to prepare a state which has arbitrarily sharp distribution of position and momentum. (There's no simultaneous measurement here.)



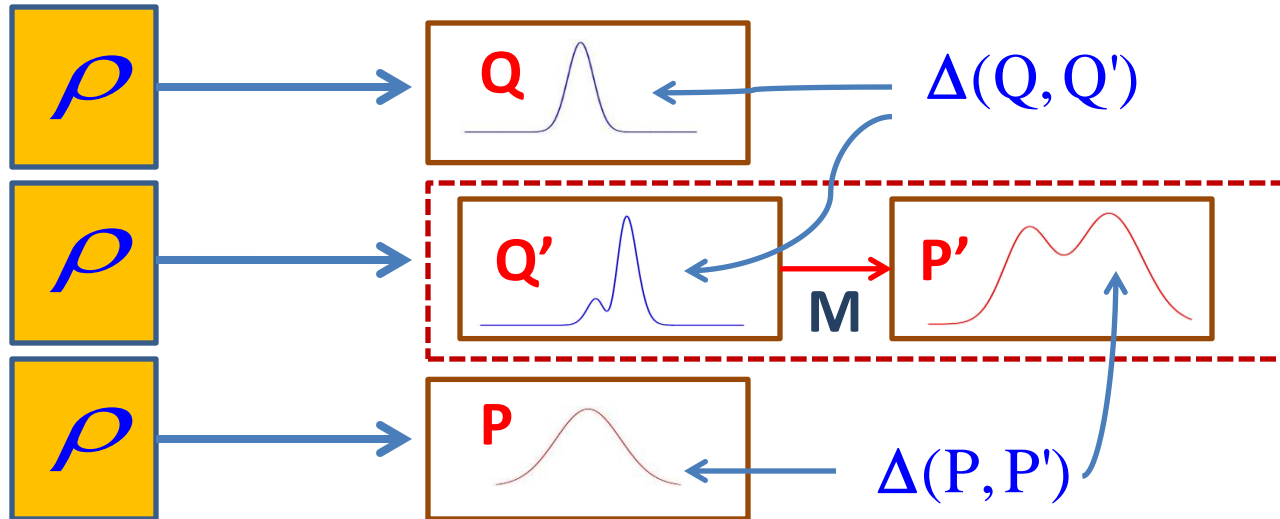
- No matter what the initial state is, we have:

$$\Delta_\rho(Q)\Delta_\rho(P) \geq \frac{\hbar}{2}$$



How to define error and disturbance?

Werner's approach



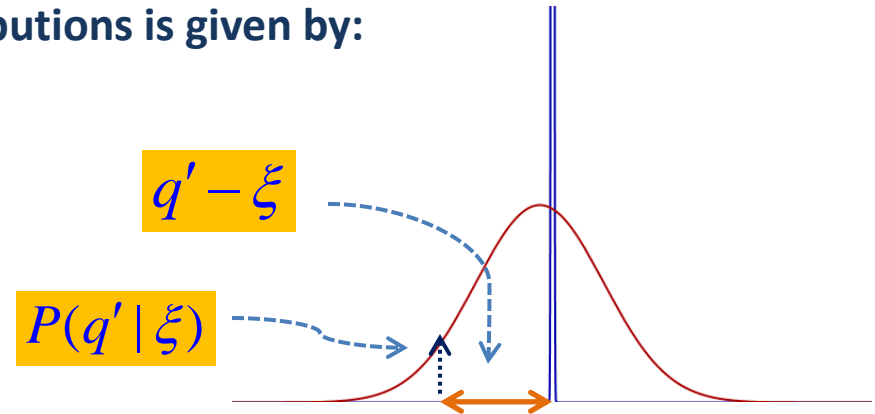
- Error in position measurement is the difference of the two distributions obtaining from Q' and Q .
- Disturbance on momentum caused by position measurement is the difference of two exact momentum measurements.

Quantifying error and disturbance

Let's consider the distributions which are sharply distributed around a position

Value ξ , the difference of the two distributions is given by:

$$D(\rho, Q'; \xi) = \left\langle (q' - \xi)^2 \right\rangle_{\rho, Q'}$$

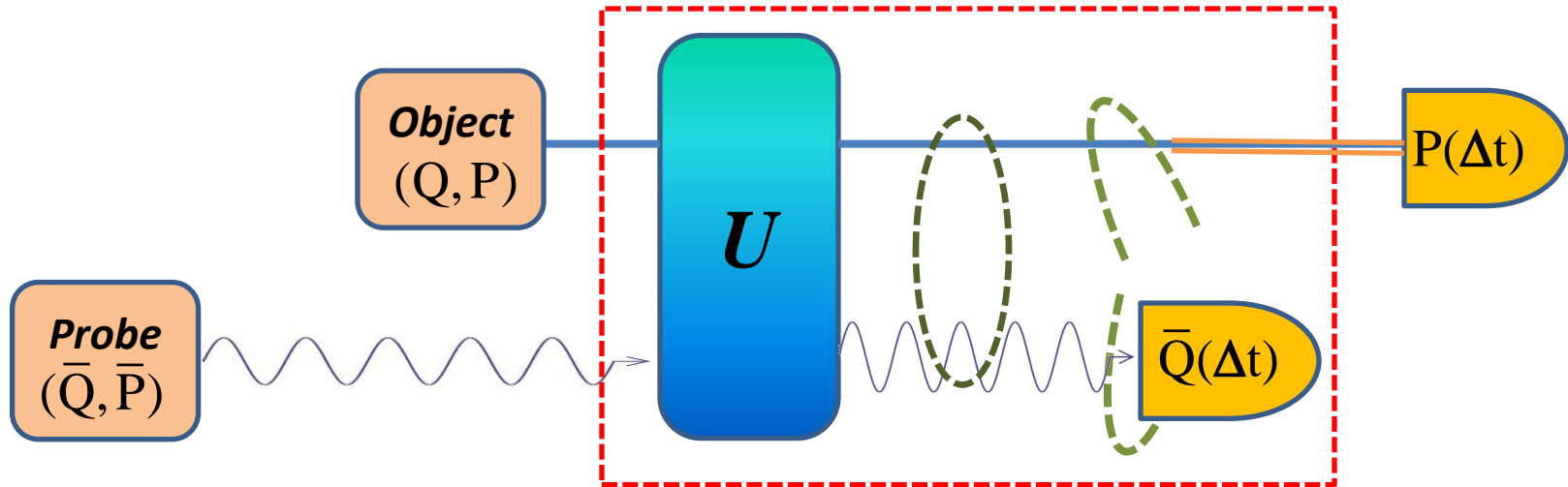


$$\Delta_c(Q, Q') = \lim_{\epsilon \rightarrow 0} \left[\sup \left\{ D(\rho, Q'; \xi) \mid \rho, \xi; D(\rho, Q; \xi) \leq \epsilon \right\} \right]$$

$$\Delta_c(P, P') = \lim_{\epsilon \rightarrow 0} \left[\sup \left\{ D(\rho', P'; \zeta) \mid \rho', \zeta; D(\rho, P; \zeta) \leq \epsilon \right\} \right]$$

$$\Delta_c(Q, Q') \Delta_c(P, P') \geq \frac{\hbar}{2}$$

Ozawa's EDR relation



- To avoid physically inaccessible resources to break Heisenberg's EDR:

At t=0, for both probe and the system wave functions.

$$\left\| \overline{Q}^m \overline{P}^n \xi \right\| < \infty, \quad \forall \{m, n\}$$

- Gauss introduced the *mean square error and disturbance* as:

$$\epsilon(Q, \psi) = \epsilon_G(Q(0), Q(\Delta t))$$

$$\epsilon_G(\Theta, \Omega) = \left(\int \int_{\mathbb{R}^2} (\omega - \theta)^2 \mu^{\Theta, \Omega}(d\theta, d\omega) \right)^{1/2}$$

- $\Theta = \theta$ is the true value and $\Omega = \omega$ is the measured value.

- **Let's suppose the interaction Hamiltonian to be:**

$$H(\alpha, \beta, \gamma) = \alpha \left(QP - \overline{QP} \right) + \beta \overline{QP} + \gamma Q\overline{P}$$

- **Solving Heisenberg's equations of motion:**

$$Q(\Delta t) = a Q(0) + b \overline{Q}(0) \quad , \quad \overline{Q}(\Delta t) = c Q(0) + d \overline{Q}(0)$$

$$P(\Delta t) = d P(0) - c \overline{P}(0) \quad , \quad \overline{P}(\Delta t) = -b P(0) + a \overline{P}(0)$$

- **The error at this case would be:**

$$\epsilon(Q, \psi) = \left\| \left[(c-1)Q(0) + d \overline{Q}(0) \right] (\psi \otimes \xi) \right\| < \infty$$

- **Root mean square disturbance:**

$$\eta(P, \psi) = \left\| \left[(d-1)P(0) - c \overline{P}(0) \right] (\psi \otimes \xi) \right\| < \infty$$

- **Error Disturbance Relation:**

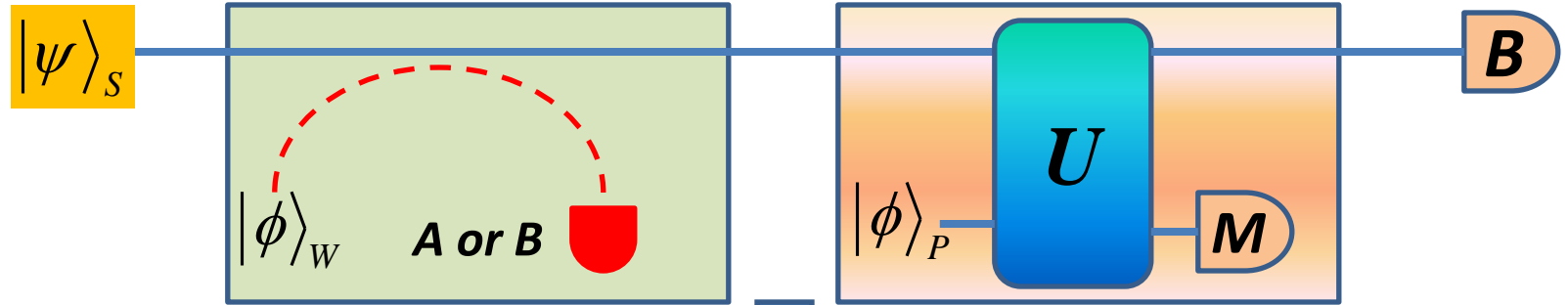
$$\epsilon(Q, \psi) \eta(P, \psi) \geq \frac{|1 - c - d| \hbar}{2}$$

- **Error free measurements:**

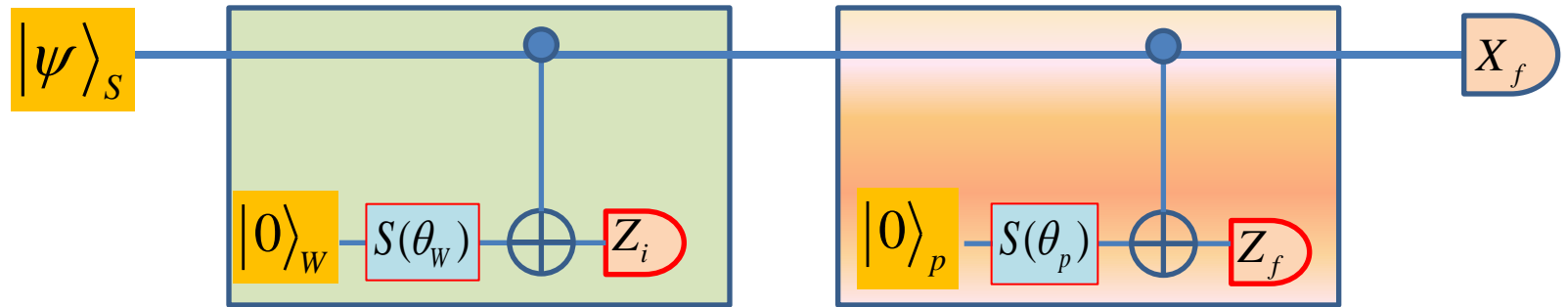
$$b = -1, c = 1, d = 0, a > -2$$

$$\begin{cases} \epsilon(Q, \psi) = 0 \\ \eta(P, \psi) < \infty \end{cases} \Rightarrow \epsilon(Q, \psi) \eta(P, \psi) = 0$$

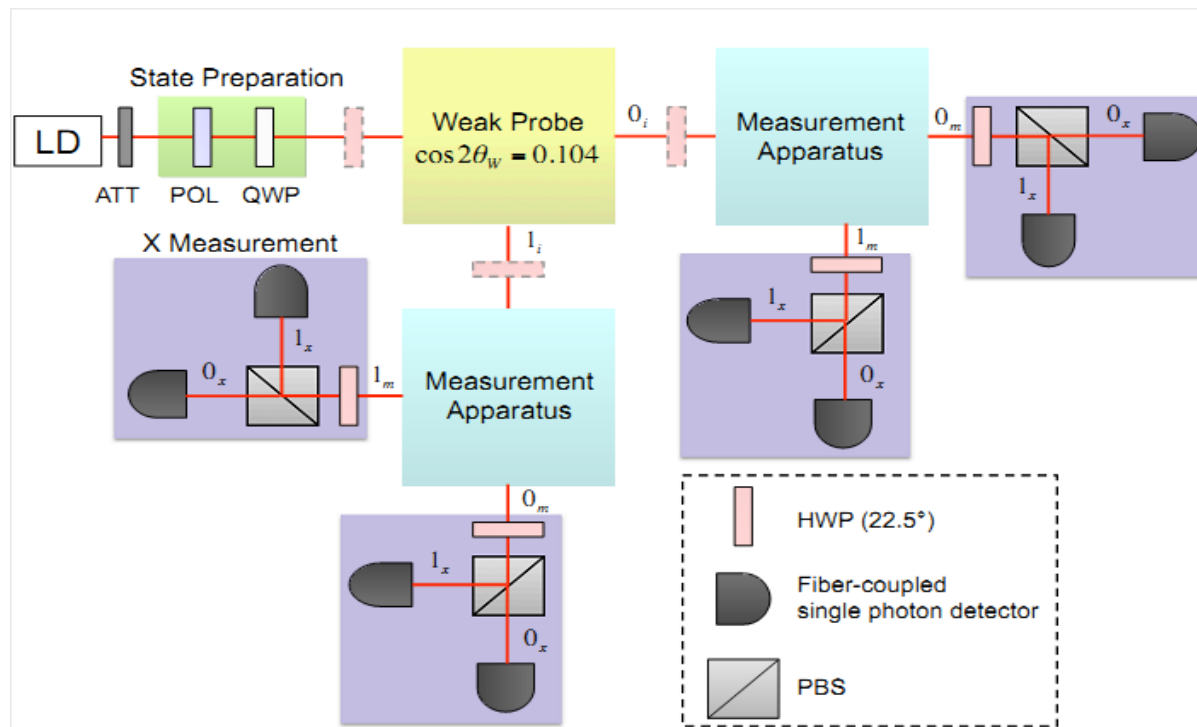
Experimental test of EDR



Spin measurement in directions X and Z:



$$S(\theta) = \begin{pmatrix} \cos \theta & \sin \theta \\ \sin \theta & -\sin \theta \end{pmatrix}, \quad E_{Z_f} = \frac{1}{2} (I \pm (\cos 2\theta) Z)$$

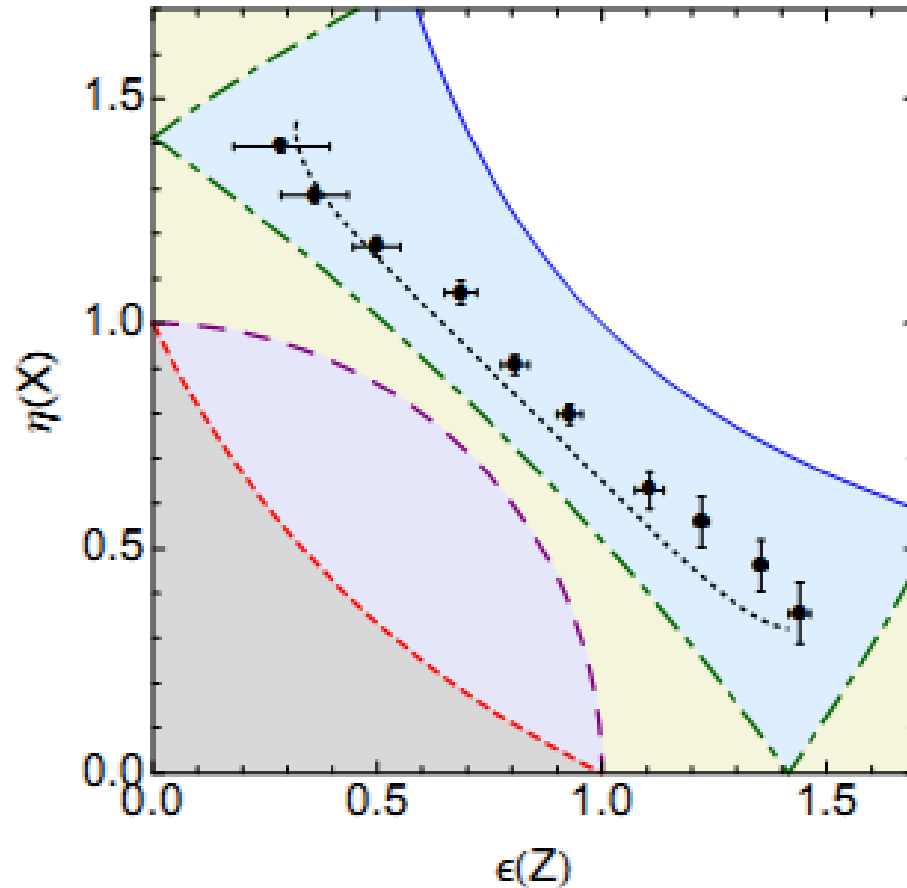


The error in measurement of Z, is given by:

$$\epsilon(Z)^2 = 2 \left(1 - \frac{1}{\cos \theta_w} \sum_{i,f} z_f z_i P(z_f, z_i) \right)$$

$$P(z_f = 1, z_i = 1) = \sum_k N_{00k} / \sum_{i,k} N_{i0k}$$

The disturbance of X, is found by substituting x_f and x_i instead of z_f and z_i .



Blue: Heisenberg's bound---Red (short dashed): Ozawa's bound---purple (long dashed) and green (dot-dashed): Branciard's bound---black dot curve: theoretical prediction---Black (filled) circles: experimental data.

Appleby's Formulation for EDR

$$\bar{\epsilon}(Q) = \sup_{\psi} \epsilon(Q, \psi) , \bar{\eta}(P) = \sup_{\psi} \eta(P, \psi)$$

- Using This relation Appleby claimed that:

$$\bar{\eta}(P)\bar{\epsilon}(Q) \geq \frac{\hbar}{2}$$

- Appleby's formulation does not work for error-free measurements!

$$\begin{cases} \bar{\epsilon}(Q) = 0 \\ \bar{\eta}(P) \rightarrow \infty \end{cases} \Rightarrow \bar{\epsilon}(Q)\bar{\eta}(P) = ?$$

- Since we can check only a limited set of ψ ; the last term is practically zero.

- The error and disturbance introduced by Appleby are almost always infinite.

$$\bar{\epsilon}(\mathbf{Q}) = \begin{cases} d \|\bar{\mathbf{Q}}\xi\| & , \quad c = 1 \\ \infty & , \quad c \neq 1 \end{cases} \quad \bar{\eta}(\mathbf{P}) = \begin{cases} d \|\bar{\mathbf{P}}\xi\| & , \quad d = 1 \\ \infty & , \quad d \neq 1 \end{cases}$$

- For linear measurements Ozawa showed that Werner's definition of error and disturbance is the same as that of Appleby's.



Thank you

Red

Blue

Purple

Orange

White