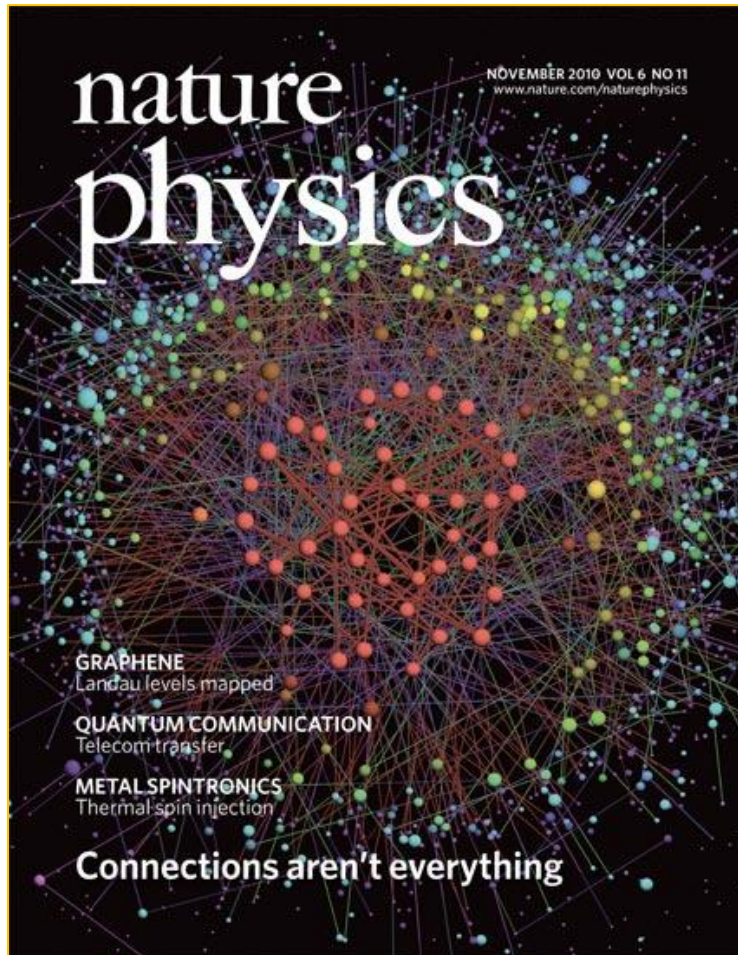


Quantum Random Networks



Zahra Raissi

Based on:

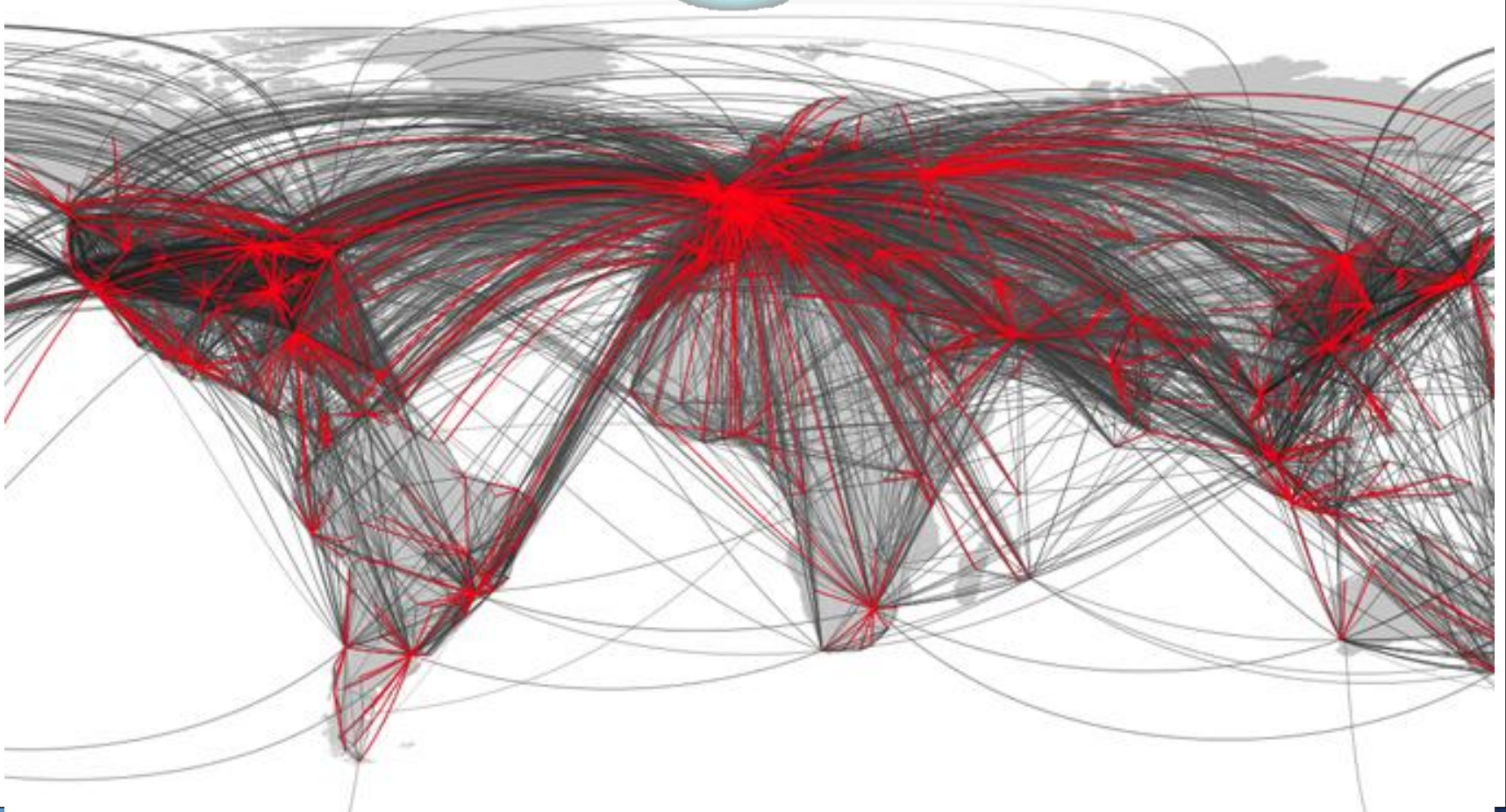
- S. Perseguers, M. Lewenstein, A. Acin, J. I. Cirac "Quantum random networks", Nature Physics 6, 539 - 543 (2010)

■ Classical Network

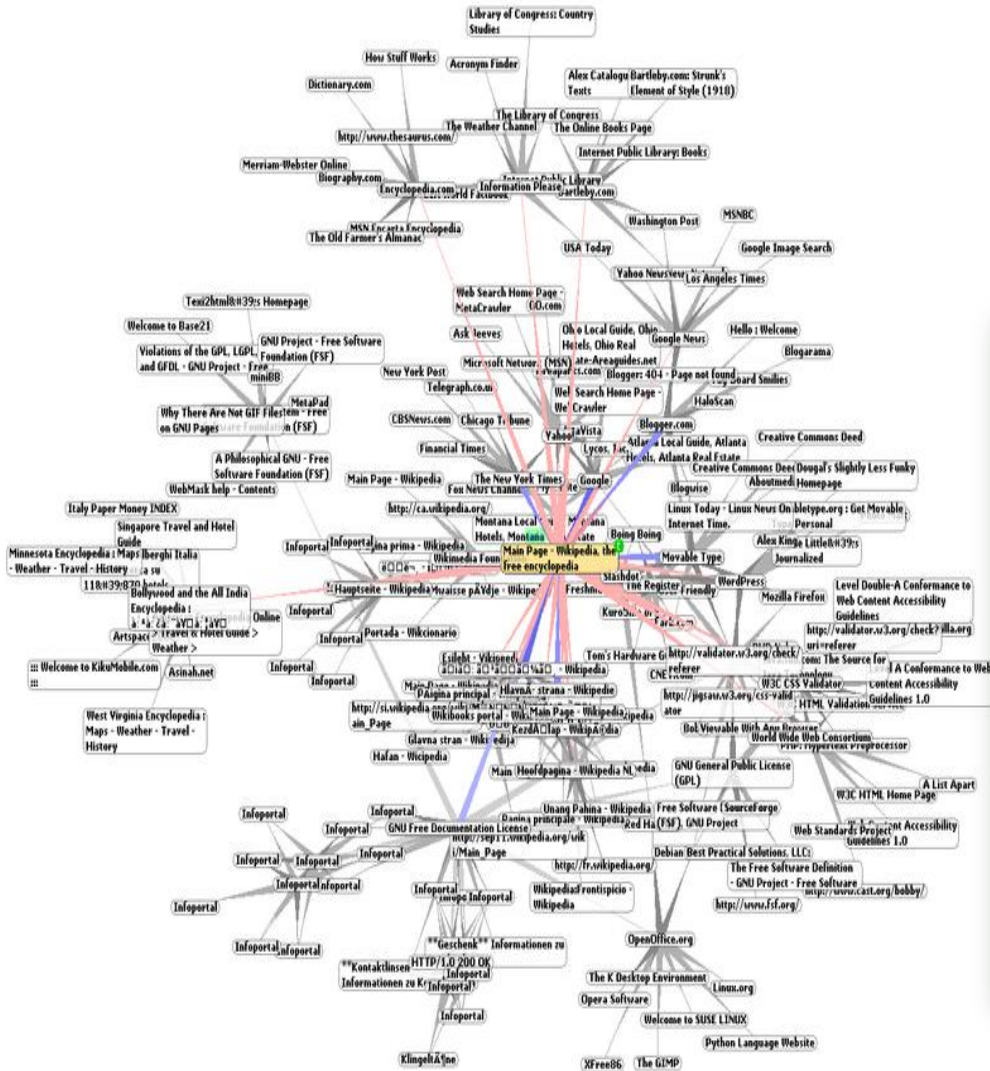
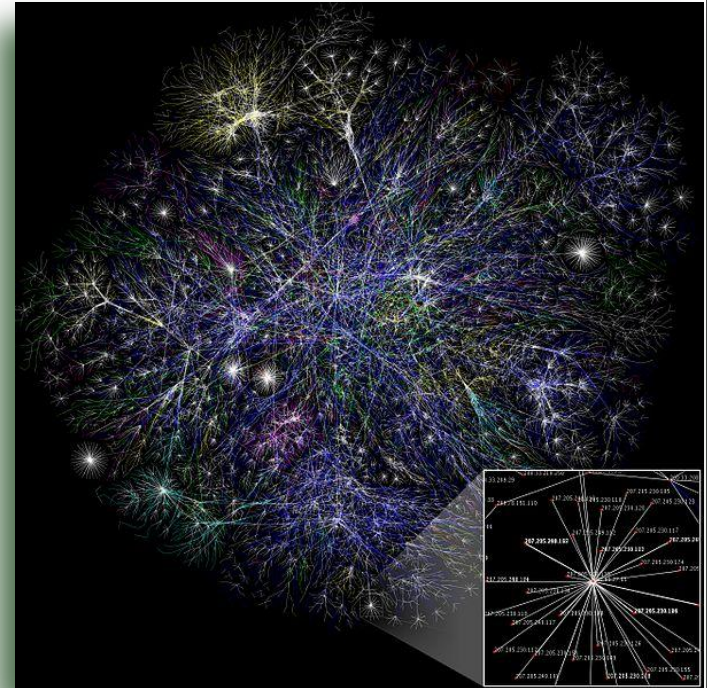
- Some Examples
- Erdős-Rényi model



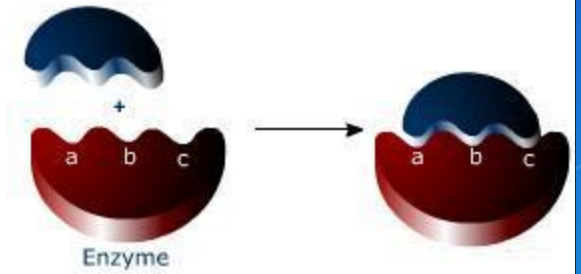
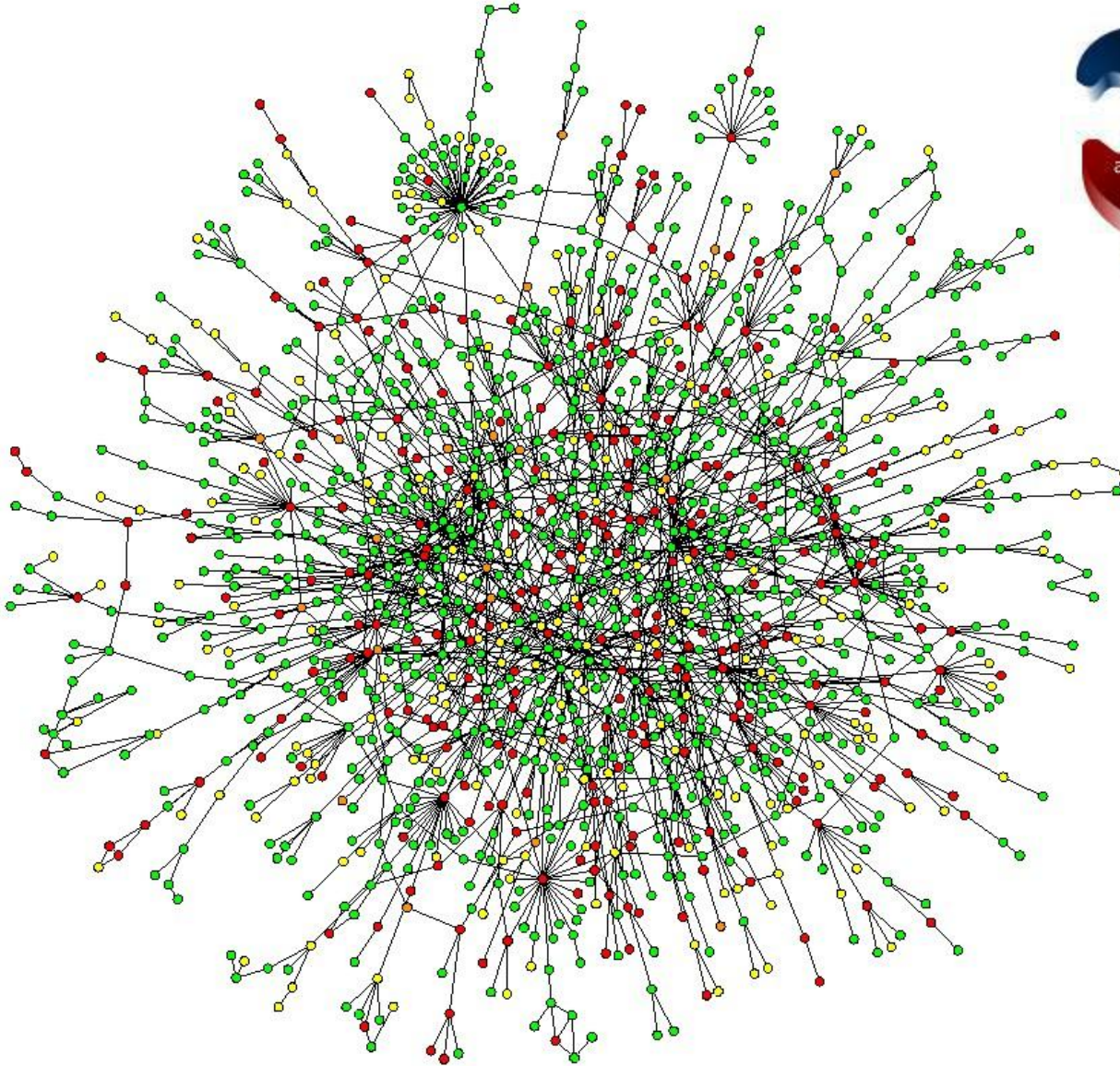
Air transportation



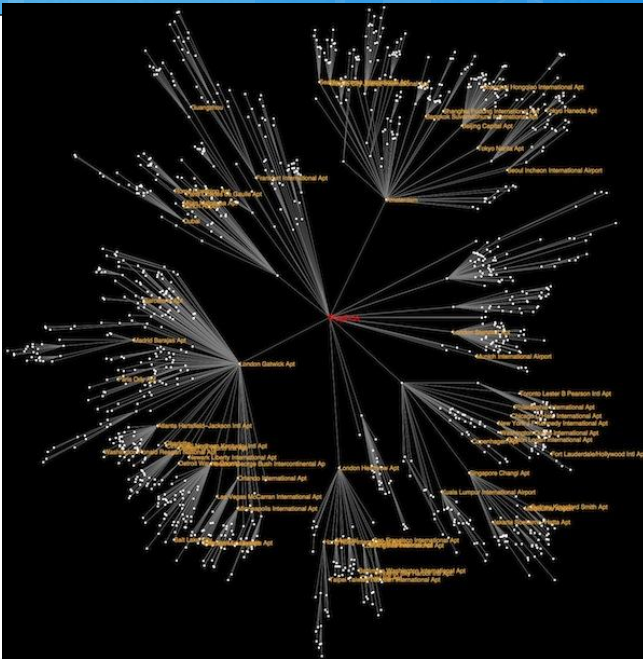
Internet



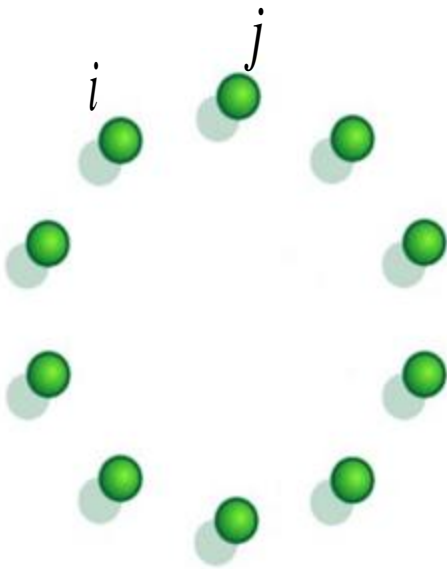
Protein network



spread of the disease



Erdős and Rényi



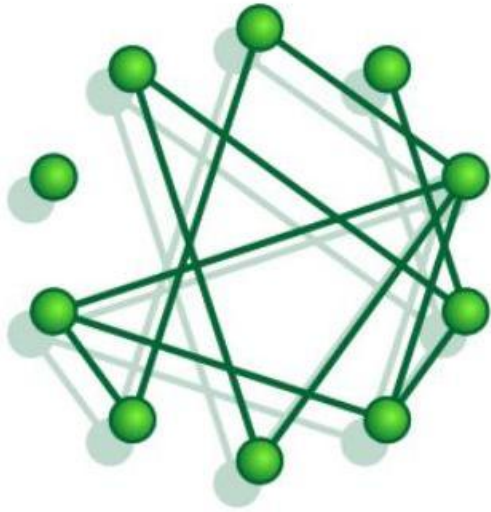
The simplest and
most studied
model



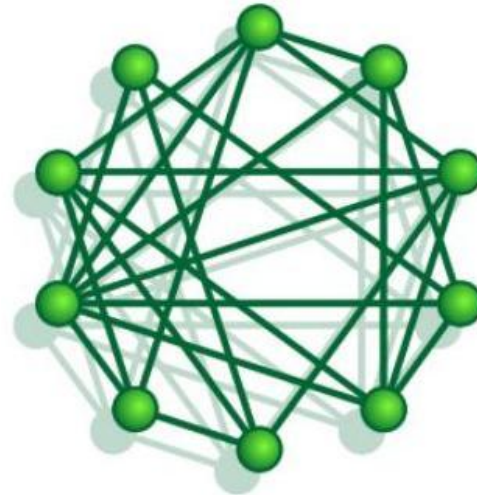
$$p_{ij}$$

$$p_{ij} = p$$

Erdős and Rényi



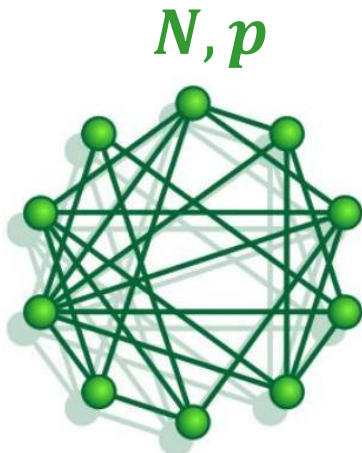
$$p = \frac{1}{6}$$



$$p = \frac{1}{2}$$

Probability in graphs

- The number of links attached to a node?




$$\text{prob}(m) = \binom{N-1}{m} p^m (1-p)^{N-1-m}$$

Select m nodes from $N-1$ → $\binom{N-1}{m}$

probability of having m edges → p^m






probability of missing $N-1-m$ edges → $(1-p)^{N-1-m}$

- The probability for the appearance a  ?



$$\text{prob}(\text{triangle}) = p^3$$






Classical Table

$p_c \sim$	$\frac{1}{N^{\frac{3}{2}}}$	$\frac{1}{N^{\frac{4}{3}}}$	$\frac{1}{N}$	$\frac{1}{N^{\frac{2}{3}}}$
			 	

if $p \geq p_c(N)$  Appearance of subgraph F with n nodes and l

$$p_c(N) = \frac{c}{N^{\frac{l}{n}}}$$




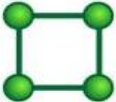

Large N

$p_c \sim$	$\frac{1}{N^{\frac{3}{2}}}$	$\frac{1}{N^{\frac{4}{3}}}$	$\frac{1}{N}$	$\frac{1}{N^{\frac{2}{3}}}$
			 	

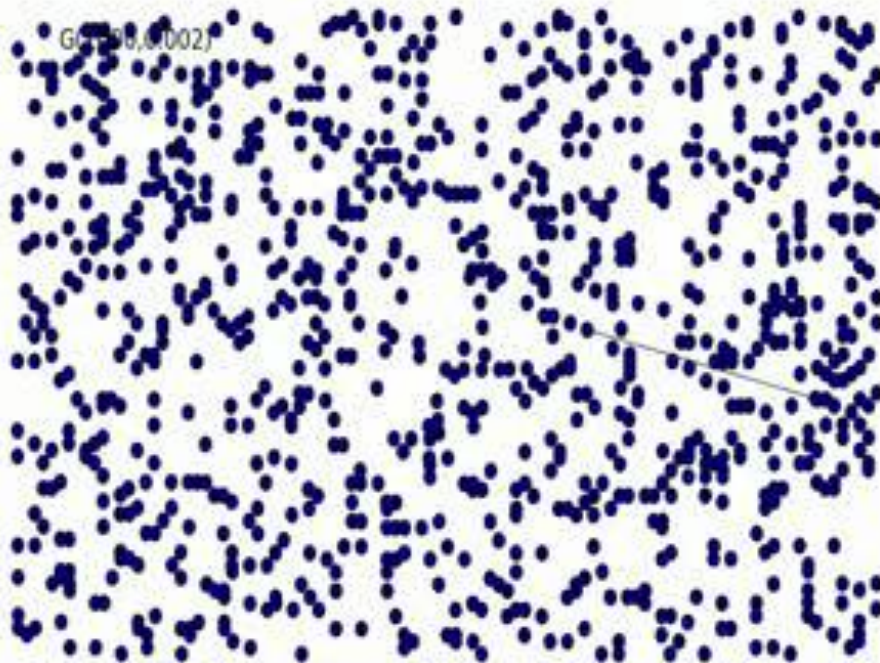
$$p \geq p_c = \frac{cte}{N^{\frac{4}{3}}}$$

$$\# \text{ nodes} = \binom{N}{4} \sim N^4 \quad \Rightarrow \quad \langle \text{star graph} \rangle = p^3(N^4) \geq 1 \Rightarrow p \geq \frac{1}{N^{\frac{4}{3}}}$$

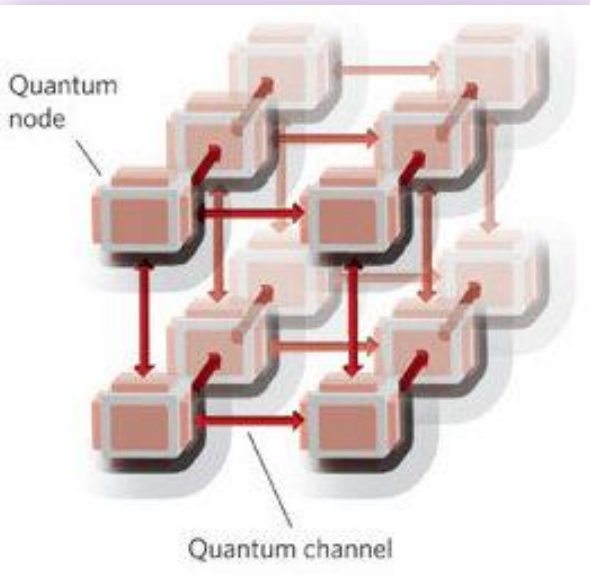
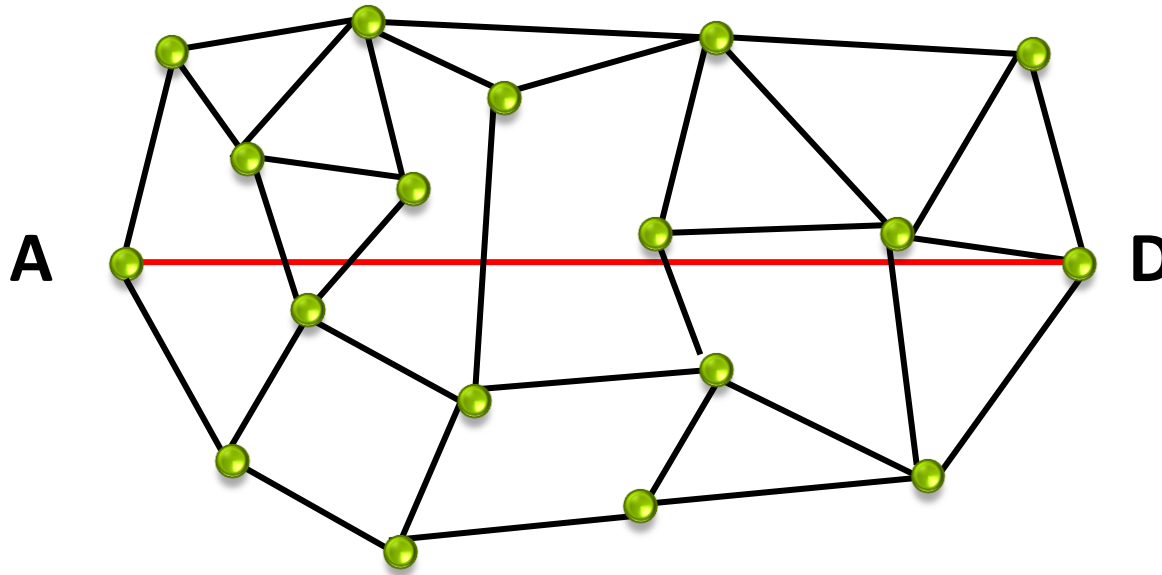
Erdős and Rényi

$p_c \sim$	$\frac{1}{N^{\frac{3}{2}}}$	$\frac{1}{N^{\frac{4}{3}}}$	$\frac{1}{N}$	$\frac{1}{N^{\frac{2}{3}}}$	
					

$G(N, p)$

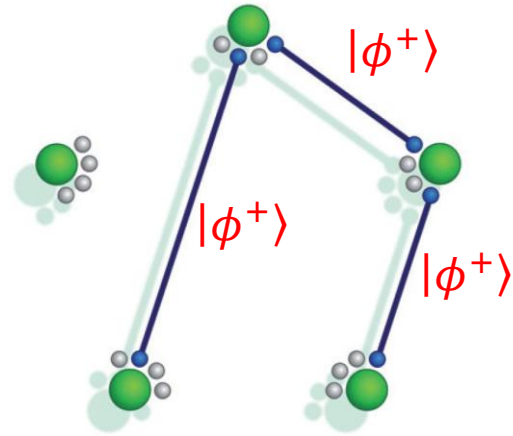
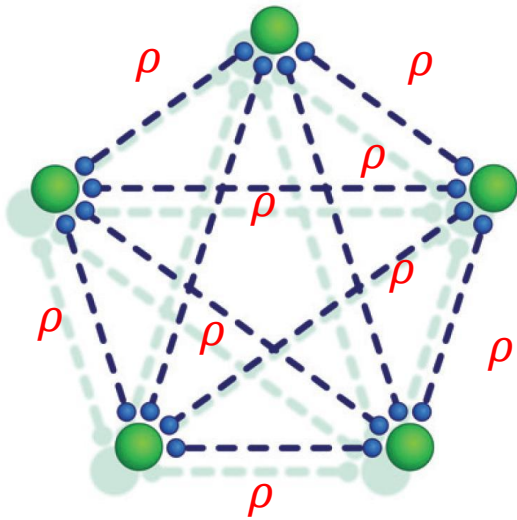


Quantum random graphs



- **Teleportation**
- **Quantum State Transfer**
- **cryptography**
- **Quantum Communication**
- **Distributed Computing**
- **Quantum Information Transfer**

Quantum random graphs



$$p_{ij} \rightarrow \rho_{ij}$$

$$p \rightarrow \rho$$

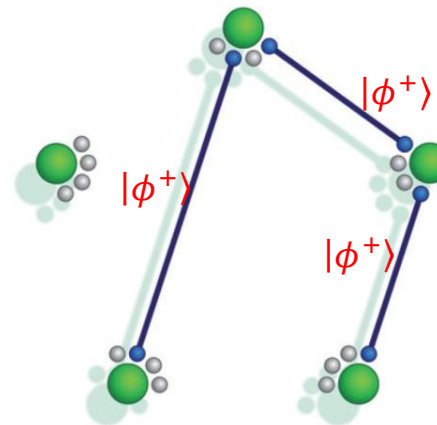
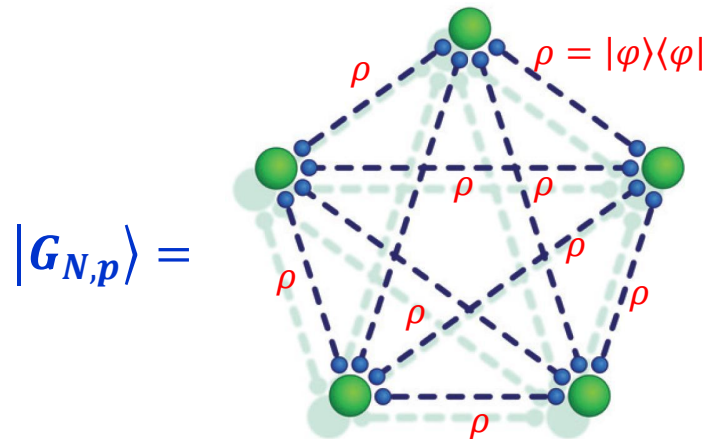
$$\rho = |\varphi\rangle\langle\varphi|$$

$$|\varphi\rangle = \sqrt{1 - \frac{p}{2}}|00\rangle + \sqrt{\frac{p}{2}}|11\rangle \text{ where } 0 \leq p \leq 1$$

Majorization theory:

$$\text{prob}(|\varphi\rangle \rightarrow |\phi^+\rangle) = 2 \left(\frac{p}{2}\right) = p$$

Main Results



$$|\varphi\rangle = \sqrt{1 - \frac{p}{2}}|00\rangle + \sqrt{\frac{p}{2}}|11\rangle$$

Majorization theory:

$$\text{prob}(|\varphi\rangle \rightarrow |\phi^+\rangle) = p$$

➤ Min probability requirement?

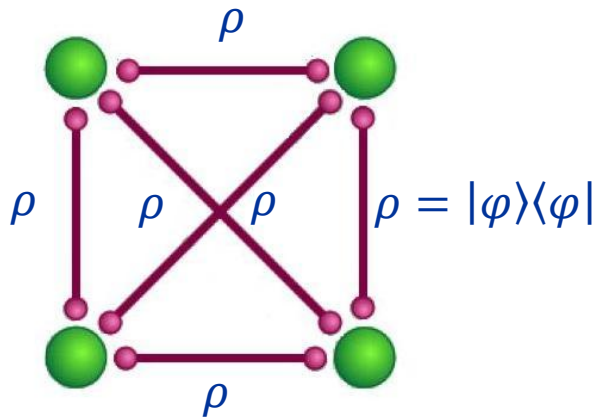
$$|\langle G_{N,p} | 0 \cdots 0 \rangle|^2$$

✓ in which case no local quantum operation is able to create entanglement between the nodes.



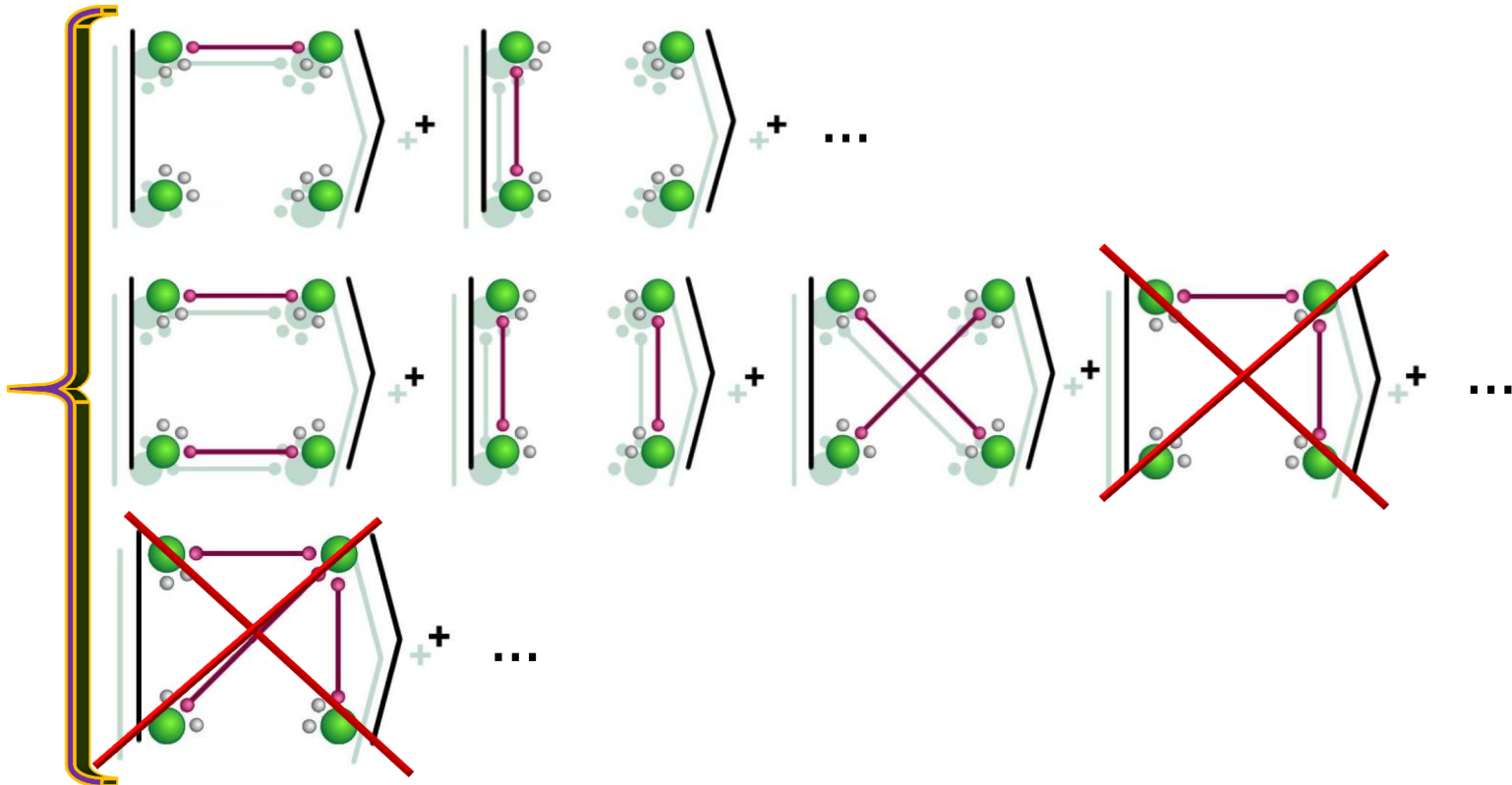
$$p = \frac{2c}{N^2}$$

GHZ states of four qutrits



$$|\varphi\rangle = \sqrt{1 - \frac{p}{2}}|00\rangle + \sqrt{\frac{p}{2}}|11\rangle, \quad p = \frac{2c}{N^2}$$

After LOCC
 we know $\text{prob}(m \geq 2) = 0$

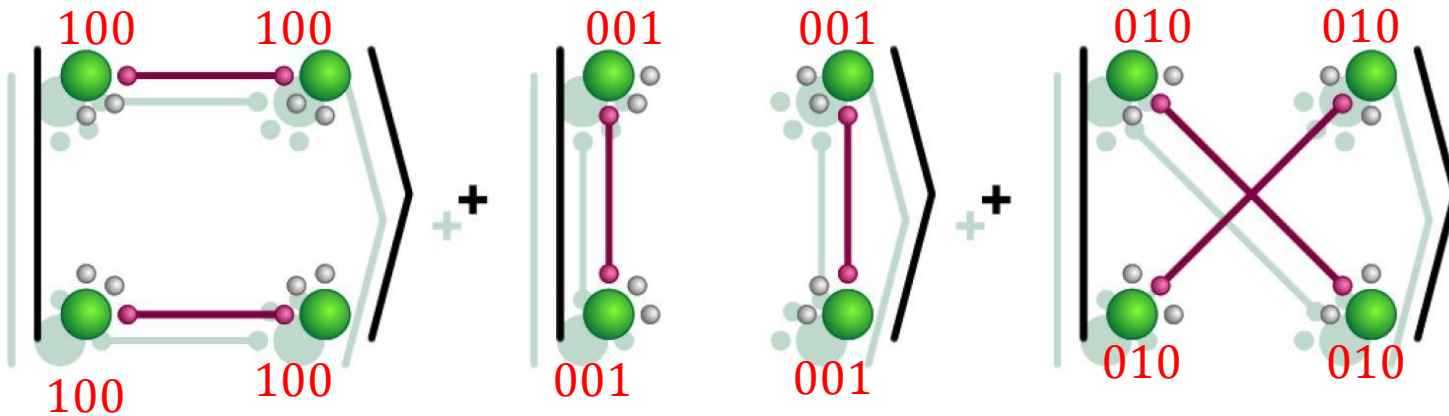


GHZ states of four qutrits

$$P_m = \sum_{\pi_m} \pi_m \underbrace{|0 \dots 0\rangle}_{N-m-1} \underbrace{|1 \dots 1\rangle}_m \langle 0 \dots 0 | \langle 1 \dots 1 | \pi_m^\dagger$$



$$m = 1$$



$$|100100100100\rangle + |001001001001\rangle + |010010010010\rangle$$

1 1 1 1
2 2 2 2
3 3 3 3

$$|GHZ\rangle \propto |1111\rangle + |2222\rangle + |3333\rangle$$

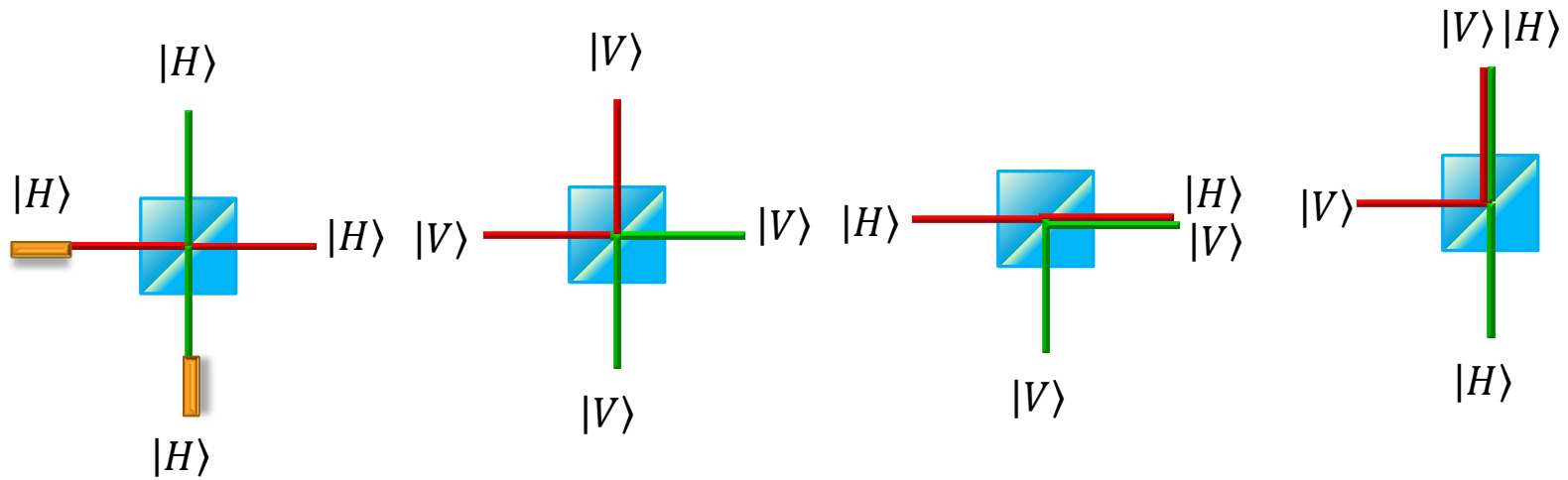
Experiment ...?

Create GHZ states

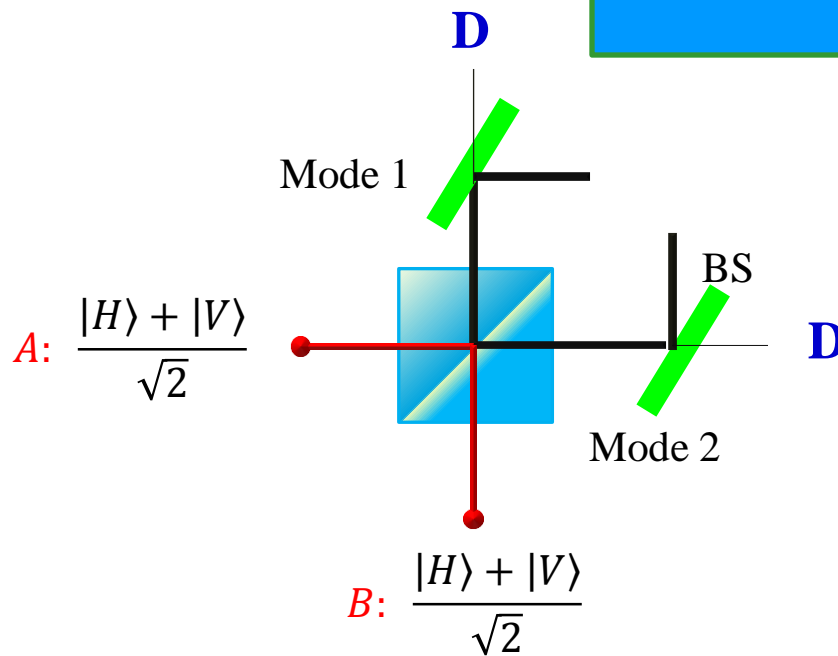


Polarizing Beam Splitter

$$\frac{|H\rangle + |V\rangle}{\sqrt{2}}$$



Bell State

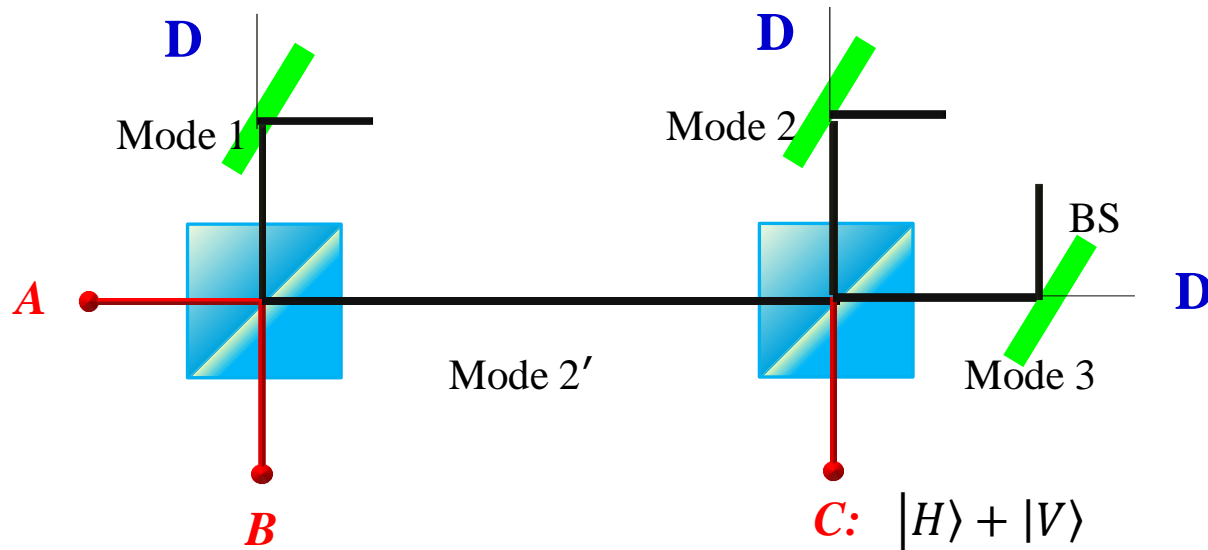


$$|\psi_{in}\rangle = \underbrace{|H_A\rangle |H_B\rangle}_{\text{HH}} + \underbrace{|H_A\rangle |V_B\rangle}_{\text{HV}} + \underbrace{|V_A\rangle |H_B\rangle}_{\text{VH}} + \underbrace{|V_A\rangle |V_B\rangle}_{\text{VV}}$$

$$|\psi_{PBS}\rangle = |H_2\rangle |H_1\rangle + |H_2\rangle |V_2\rangle + |V_1\rangle |H_1\rangle + |V_1\rangle |V_2\rangle$$

$$|\psi_{final}\rangle = |H_1\rangle |H_2\rangle + |V_1\rangle |V_2\rangle = |\phi^+\rangle$$

Create GHZ states



$$|\phi^+\rangle = \underbrace{|H_1\rangle |H_2\rangle}_I + \underbrace{|V_1\rangle |V_2\rangle}_II$$



$$I: |H_1\rangle |H_3\rangle |H_2\rangle + |H_1\rangle |H_3\rangle |V_3\rangle$$

$$II: |V_1\rangle |V_2\rangle |H_2\rangle + |V_1\rangle |V_2\rangle |V_3\rangle$$

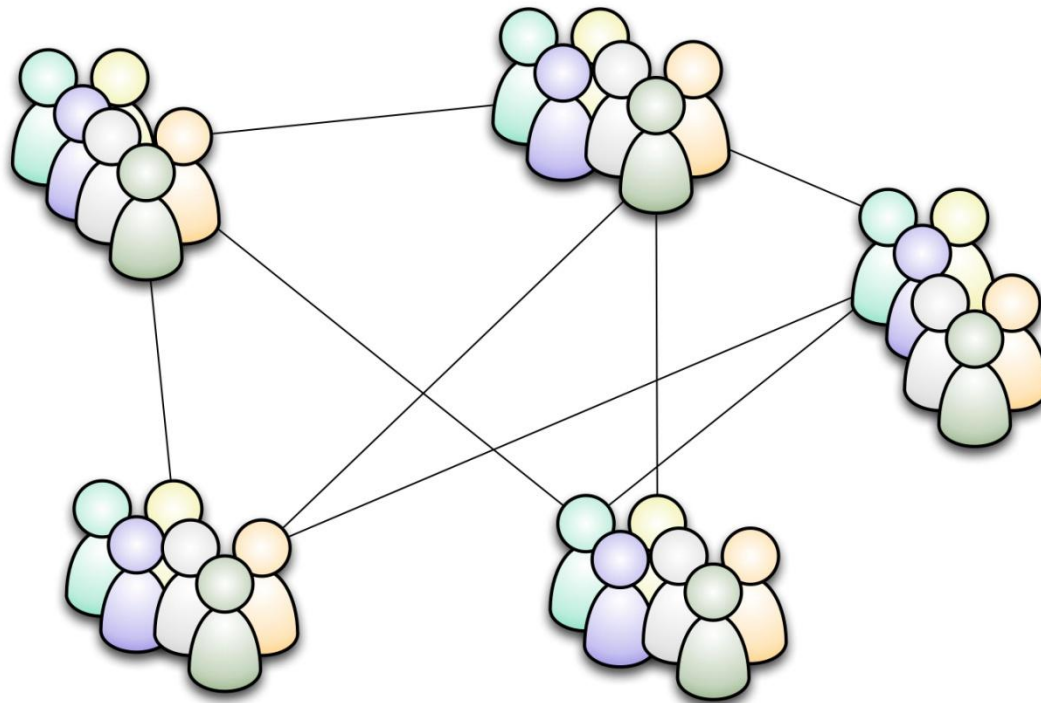
$$|GHZ\rangle = |H\rangle |H\rangle |H\rangle + |V\rangle |V\rangle |V\rangle$$

Our Problem

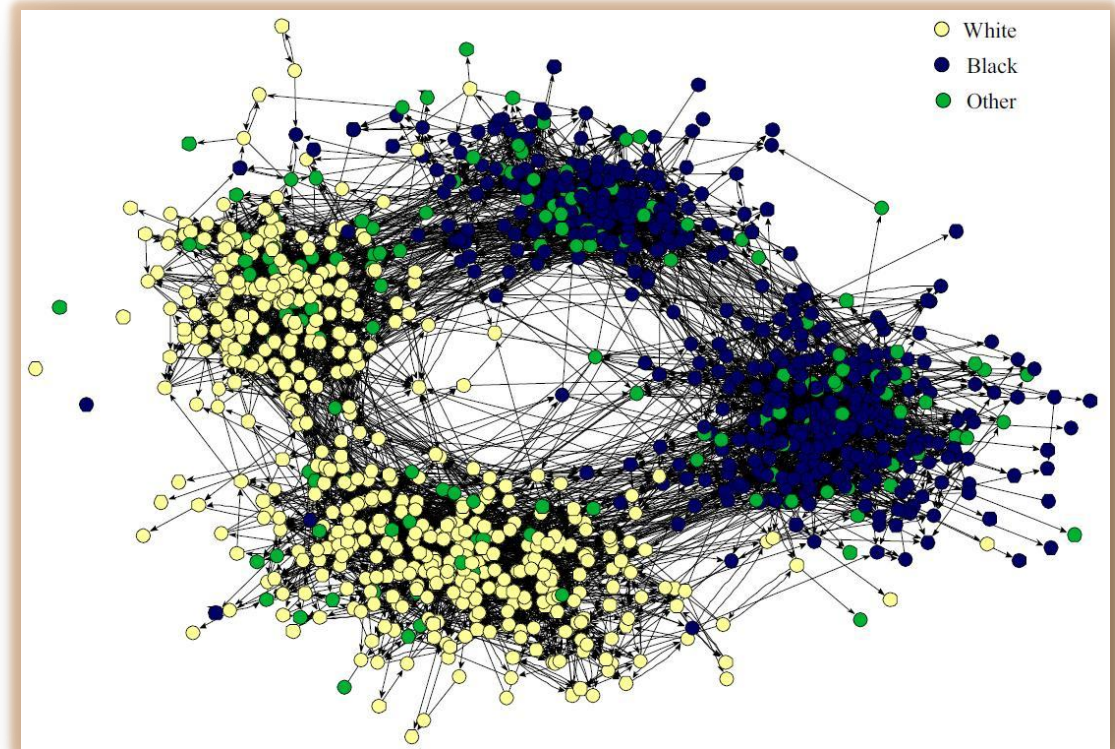
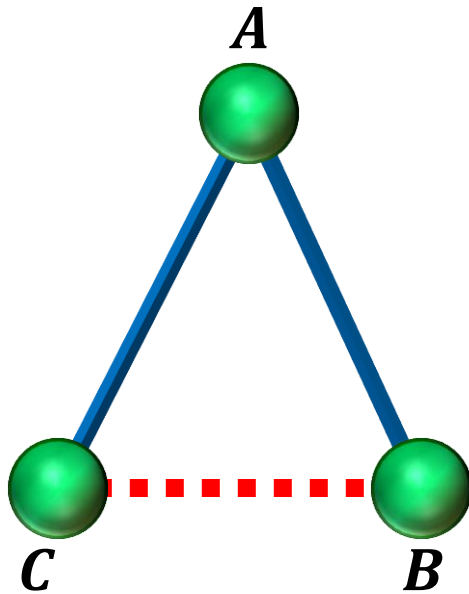


Small World Property

- shortest paths between nodes $\sim \ln(N)$, where N is the number of nodes.

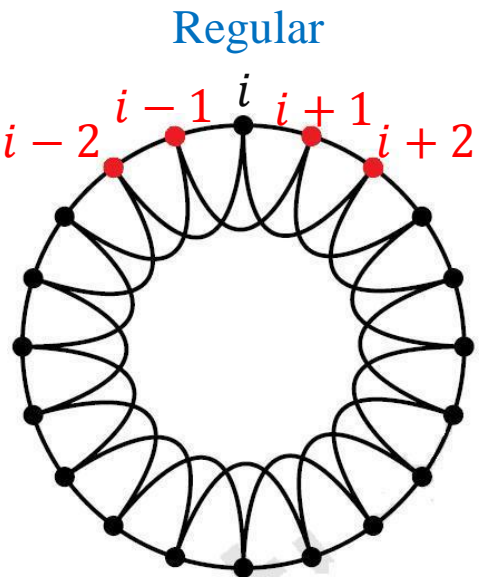


Clustering

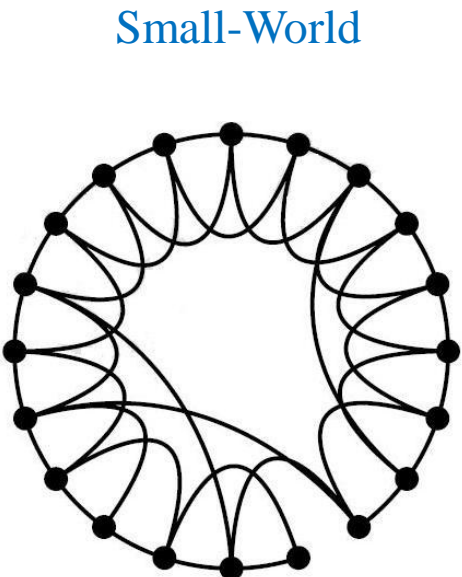


✓ These properties distinguish real networks from simple models such as the Erdős Renyi random graph model

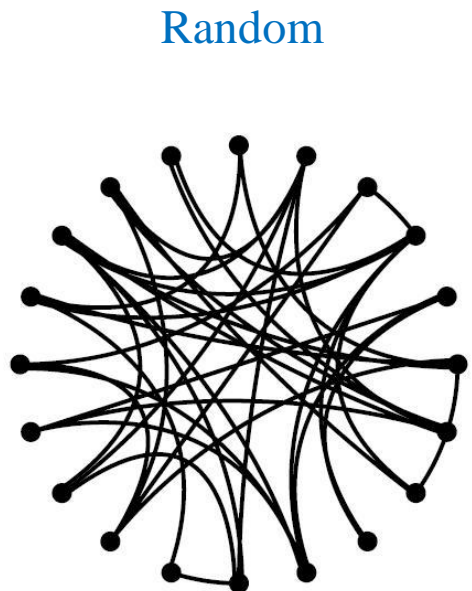
Real Networks



$p = 0$



increasing randomness



$p = 1$

- high clustering
- high average path length
 $\bar{l} \propto N$

- high clustering
- small average path length
 $\bar{l} \propto \ln(N)$

- small clustering
- small average path length
 $\bar{l} \propto \ln(N)$

D. J. Watts and S. H. Strogatz, *Nature*, **393**, 440 (1998).



Thanks for your attention.