

18.09.06 - 19.09.11

Also: $\mathbb{RP}^2 = \text{Real projective plane.}$

$$\mathbb{RP}^2 = \frac{\mathbb{R}^3}{v \sim \lambda v} \rightarrow \text{Coordinate charts: } \mathbb{R}^3 = (x^1, x^2, x^3)$$

$$U_1: x_1 \neq 0$$



$$\xi_{(1)}^2 = \frac{x^2}{x^1}$$

$$\xi_{(1)}^3 = \frac{x^3}{x^1}$$

$$\alpha = \frac{x^2}{x^1}$$

$$\beta = \frac{x^3}{x^1}$$

$$U_2: x_2 \neq 0$$

$$\xi_{(2)}^1 = \frac{x^1}{x^2}$$

$$\xi_{(2)}^3 = \frac{x^3}{x^2}$$

$$\alpha' = \frac{x^1}{x^2}$$

$$\beta' = \frac{x^3}{x^2}$$

$$U_3: x_3 \neq 0$$

$$\xi_{(3)}^1 = \frac{x^1}{x^3}$$

$$\xi_{(3)}^2 = \frac{x^2}{x^3}$$

$$\alpha'' = \frac{x^1}{x^3}$$

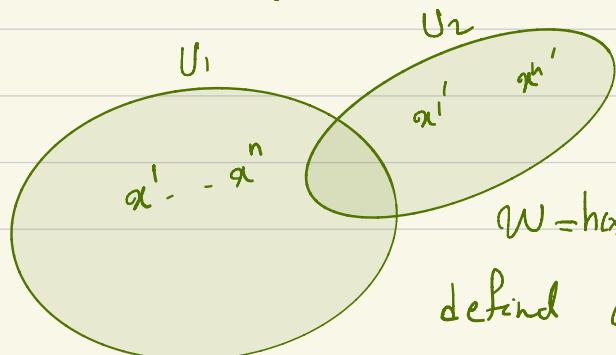
$$\beta'' = \frac{x^2}{x^3}$$

$$U_1 \cap U_2: \alpha' = \frac{1}{\alpha}, \beta' = \frac{\beta}{\alpha} \Rightarrow J = \begin{vmatrix} \frac{\partial \alpha'}{\partial x^1} & \frac{\partial \alpha'}{\partial x^2} \\ \frac{\partial \beta'}{\partial x^1} & \frac{\partial \beta'}{\partial x^2} \end{vmatrix}$$

$$= \begin{vmatrix} -\frac{1}{\alpha^2} & 0 \\ -\frac{\beta}{\alpha^2} & \frac{1}{\alpha} \end{vmatrix} = -\frac{1}{\alpha^3} \neq 0 \rightarrow \mathbb{RP}^2 \text{ is not orientable.}$$

Definition: Let M be an orientable manifold of $\dim=n$.

We can define an n -form w | $w \neq 0$ over all M .



Volume form.

$$w = h(x) dx^1 \wedge dx^2 \wedge \dots \wedge dx^n$$

$$h(x) > 0: \hookrightarrow$$

defined on U_1

$$\text{on } V_1 \cap U_2 \quad \omega = h(x(x')) \frac{\partial x^1}{\partial x^{i_1}} \frac{\partial x^2}{\partial x^{i_2}} \cdots \frac{\partial x^n}{\partial x^{i_n}} dx^{i_1} \wedge dx^{i_2} \wedge \cdots \wedge dx^{i_n}$$

$$dx^{i_1} \wedge dx^{i_2} \wedge \cdots \wedge dx^{i_n} = \epsilon^{i_1 i_2 \cdots i_n} dx^1 \wedge dx^2 \wedge \cdots \wedge dx^n \leftarrow$$

$$\text{on } V_1 \cap U_2 \quad \omega = h(x(x')) \underbrace{\frac{\partial x^1}{\partial x^{i_1}} \frac{\partial x^2}{\partial x^{i_2}} \cdots \frac{\partial x^n}{\partial x^{i_n}}} \underbrace{\epsilon^{i_1 i_2 \cdots i_n} dx^1 \wedge dx^2 \wedge \cdots \wedge dx^n}$$

$$\omega = \underbrace{h(x')} \underbrace{\|\frac{\partial x}{\partial x'}\|} dx^1 \wedge dx^2 \wedge \cdots \wedge dx^n > 0$$

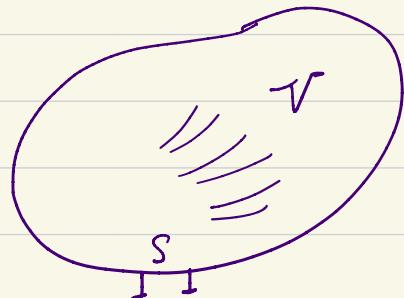
then we extend this to U_2 and then we continue on all M .

ω is positive over all M . It is called a volume form.

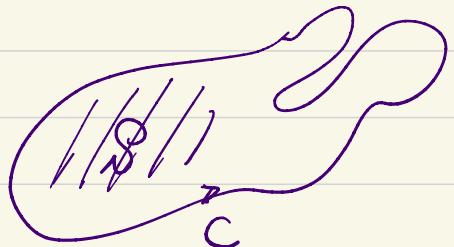
$$\text{EXAMPLE: over } S^2 \quad \omega = \sin \theta d\theta \wedge d\varphi \quad 0 < \theta < \pi$$

$\omega > 0$ over all $S^2 - \{\text{North pole}\} - \{\text{South pole}\}$.

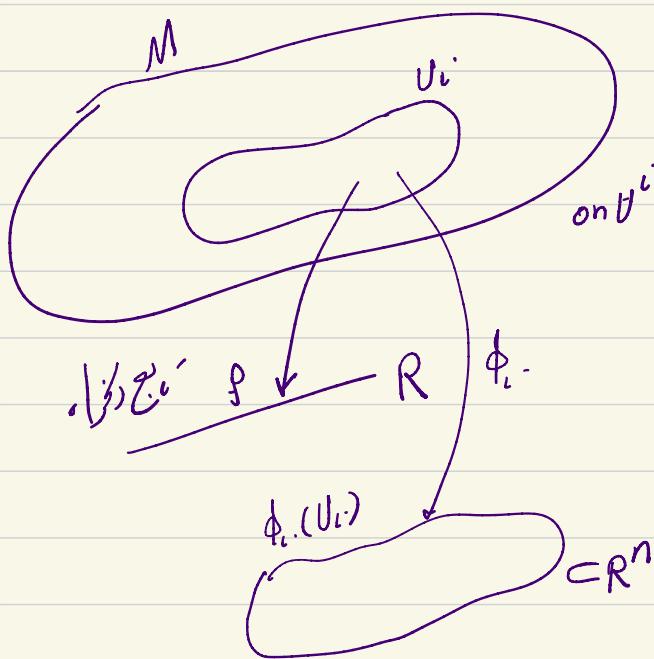
$$\text{область. } \int\limits_{\partial V} A \cdot ds = \int\limits_V (\nabla \cdot A) dV$$



$$\text{область } \int\limits_{\partial S} A \cdot d\ell = \int\limits_S (\nabla \times A) \cdot ds$$



فیض ریاضی میں اسی طرز کا انتگرال اسی طرز پر فرمائی جائے گا



تو نہیں زیر دفعہ اس سمت ہے

$$\omega = h(x) dx^n \wedge dn^n$$

$$\int f \omega := \int_{U_i} \int_{\phi_i(U_i)} f(x) h(x) dx^n \wedge dn^n - dn^n.$$

\rightarrow اسی طرز کا انتگرال

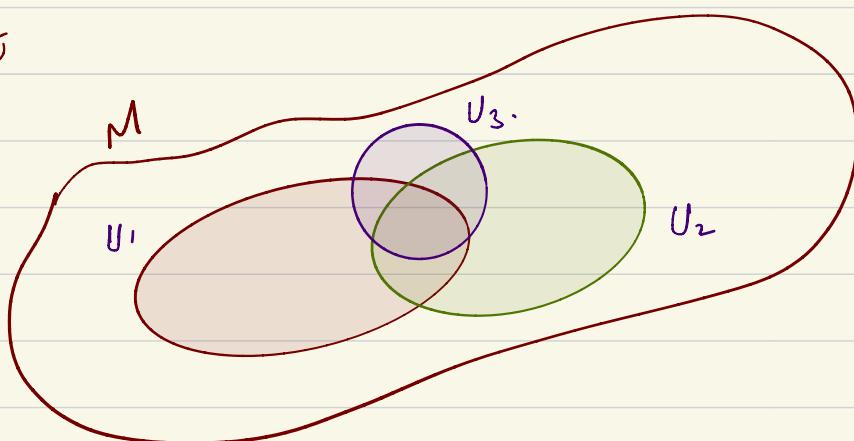
$\int_M f \omega$:

$$\int_M f \omega := \iint_{S^2} f(\theta, \varphi) \sin \theta d\theta d\varphi,$$

$\theta \in [0, \pi]$
 $\varphi \in [0, 2\pi]$

No wedge product

مودل کیلی



$$\int_M f \omega = \int_{U_1} f \omega + \int_{U_2} f \omega + \dots + \int_{U_N} f \omega$$

فیض ریاضی پر $\sum f \omega$

↓

Intuition: M is paracompact if any point of M can be covered by a finite number of charts

Defⁿ Partition of Unity:

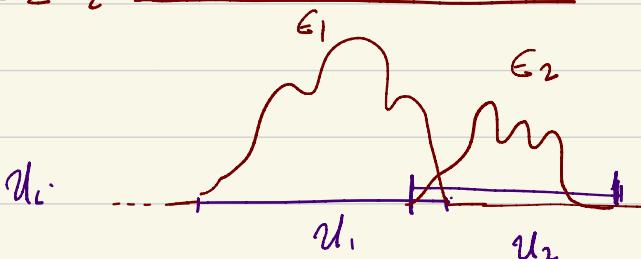
$\{\epsilon_i\}: M \rightarrow \mathbb{R}$ are continuous functions on M .

$$1 = \sum \epsilon_i$$

$$(i) \quad 0 \leq \epsilon_i(p) \leq 1$$

$$(ii) \quad \epsilon_i(p) = 0 \quad \forall p \notin U_i$$

$$(iii) \quad \sum_i \epsilon_i(p) = 1 \quad \forall p \in M$$

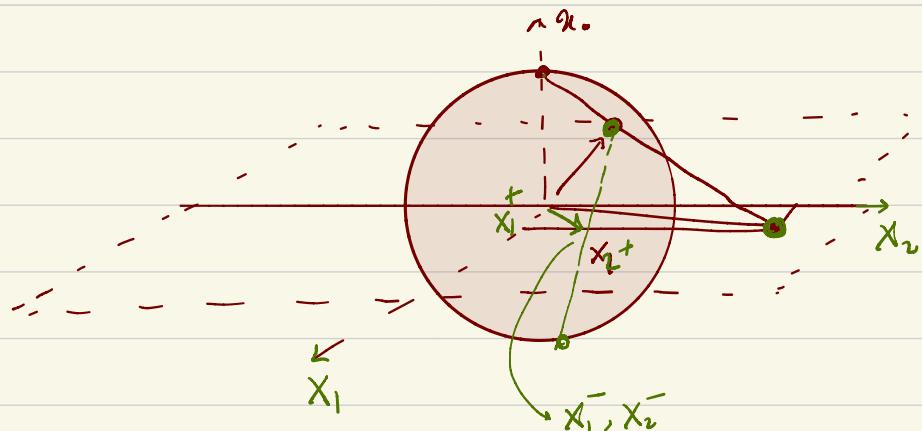


$$\int_M f \omega = \int_M f \omega (\sum_i \epsilon_i) = \sum_i \int_{U_i} f \omega \epsilon_i$$

Ex: $M = S^2$:

$$U_N = S^2 - \{N\}$$

$$U_S = S^2 - \{S\}$$



$$U_N: \quad X_1^+ = \frac{\alpha_1}{1-\alpha_0} \quad X_2^+ = \frac{\alpha_2}{1-\alpha_0}$$

$$U_S : \quad \bar{x}_1 = \frac{x_1}{1+x_0} \quad \bar{x}_2 = \frac{x_2}{1+x_0}$$

$$\begin{cases} x_0 = \cos \theta \\ x_1 = \sin \theta \cos \varphi \\ x_2 = \sin \theta \sin \varphi \end{cases} \quad \begin{aligned} x_1^+ &= \frac{\sin \theta \cos \varphi}{1 - \cos \theta} = \frac{\sin \theta \cos \varphi}{2 \sin^2 \frac{\theta}{2}} = \cot \frac{\theta}{2} \cos \varphi \\ x_2^+ &= \cot \frac{\theta}{2} \sin \varphi \end{aligned}$$

$$\text{in } U_N (\theta \neq 0) \quad x_1^+ = \cot \frac{\theta}{2} \cos \varphi \quad x_2^+ = \cot \frac{\theta}{2} \sin \varphi$$

$$Z^+ := x_1^+ + i x_2^+ = \cot \frac{\theta}{2} e^{i\varphi}$$

$$\text{in } U_S (\theta \neq \pi) \quad \bar{x}_1 = \tan \frac{\theta}{2} \cos \varphi \quad \bar{x}_2 = \tan \frac{\theta}{2} \sin \varphi$$

$$\bar{Z} = \bar{x}_1 + i \bar{x}_2 = \tan \frac{\theta}{2} e^{i\varphi}$$

$$\text{in } U_N \cap U_S : \quad Z^+ = \frac{1}{Z^-*}$$

on S^2 : we had $\omega = \sin \theta d\theta \wedge d\varphi$

$$\text{on } U_N : \quad \omega = \sin \theta \left(\frac{\partial \theta}{\partial x_1^+} dx_1^+ + \frac{\partial \theta}{\partial x_2^+} dx_2^+ \right) \wedge \left(\frac{\partial \varphi}{\partial x_1^+} dx_1^+ + \frac{\partial \varphi}{\partial x_2^+} dx_2^+ \right)$$

$$\omega = \sin \theta \begin{vmatrix} \frac{\partial \theta}{\partial x_1^+} & \frac{\partial \theta}{\partial x_2^+} \\ \frac{\partial \varphi}{\partial x_1^+} & \frac{\partial \varphi}{\partial x_2^+} \end{vmatrix} dx_1^+ \wedge dx_2^+$$

$$\omega = \sin \theta \frac{1}{|\frac{\partial X_1, X_2}{\partial \theta, \varphi}|} dX_1^+ \wedge dX_2^+$$

$\int \omega =$

$$X_1^+ = \cot \frac{\theta}{2} \cos \varphi \quad X_2^+ = \cot \frac{\theta}{2} \sin \varphi$$

$$J = \begin{vmatrix} \frac{\partial X_1^+}{\partial \theta} & \frac{\partial X_1^+}{\partial \varphi} \\ \frac{\partial X_2^+}{\partial \theta} & \frac{\partial X_2^+}{\partial \varphi} \end{vmatrix} = \begin{vmatrix} -\frac{1}{2} \frac{1}{\sin^2 \frac{\theta}{2}} \cos \varphi & -\cot \frac{\theta}{2} \sin \varphi \\ -\frac{1}{2} \frac{1}{\sin^2 \frac{\theta}{2}} \sin \varphi & \cot \frac{\theta}{2} \cos \varphi \end{vmatrix} = +\frac{1}{2} \left\{ \frac{1}{\sin^2 \frac{\theta}{2}} \right\} \cot \frac{\theta}{2} \left\{ -1 \right\}$$

$$J = -\frac{1}{2} \frac{\cos \theta/2}{\sin^3 \theta/2} \Rightarrow \omega = -2 \sin \theta \frac{\sin^3 \theta/2}{\cos \theta/2} dX_1^+ \wedge dX_2^+$$

$$\omega = -4 \left(\sin \frac{\theta}{2} \right)^4 dX_1^+ \wedge dX_2^+ \quad (X_1^+)^2 + (X_2^+)^2 = \cot^2 \frac{\theta}{2}$$

$$1 + \cot^2 \frac{\theta}{2} = \frac{1}{\sin^2 \frac{\theta}{2}} \rightarrow \left(\omega = -4 \frac{1}{(1 + (X_1^+)^2 + (X_2^+)^2)^2} dX_1^+ \wedge dX_2^+ \right)$$

$$\omega_+ = \omega \text{ on } U_N.$$

$Z \bar{Z} \wedge dZ \wedge \bar{d}\bar{Z}$ is called current

$$Z = X_1 + iX_2 \quad \bar{Z} = X_1 - iX_2 \quad dZ \wedge d\bar{Z} = -2i dX_1 \wedge dX_2.$$

$$Z \bar{Z} = X_1^2 + X_2^2 \rightarrow$$

$$\omega_+ = \frac{2i}{(1 + Z \bar{Z})^2} d\bar{Z} \wedge dZ$$

ω_+ is always non-zero & well defined on U_N .

to find $\bar{\omega}$, we use change of coordinates:

For simplicity: $z_- = \xi$ $\xi = \frac{1}{\bar{z}}$ $z = \frac{1}{\xi}$, $\bar{z} = \frac{1}{\xi}$.

$$\omega_- = \frac{2\iota}{\left(1 + \frac{1}{\xi\xi}\right)^2} \frac{-1}{\xi^2} d\xi \wedge \frac{-1}{\bar{\xi}^2} d\bar{\xi} = \frac{2\iota \cdot d\xi \wedge d\bar{\xi}}{(1 + \xi\bar{\xi})^2}$$

on U_ξ :

$$\omega_- = \frac{2\iota}{(1 + \xi\bar{\xi})^2} d\xi \wedge d\bar{\xi}$$

