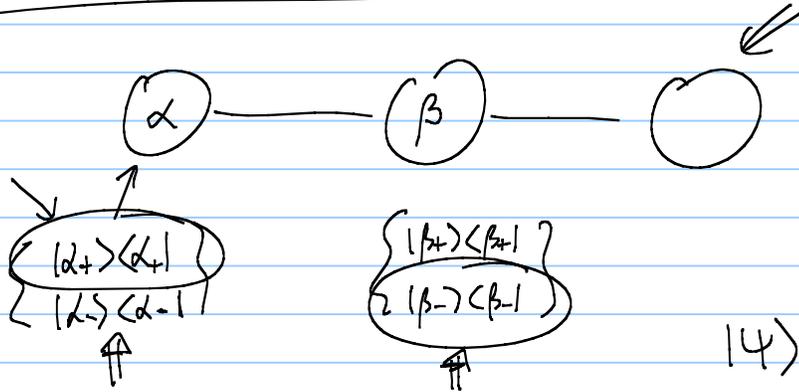


# Lecture 2



$$P_0 = |A|^2$$

$$P_1 = |B|^2$$

$$|\psi\rangle = \begin{pmatrix} A \\ B \end{pmatrix} \leftarrow \text{circuit}$$

$$|\psi'\rangle = \begin{pmatrix} B \\ A \end{pmatrix} = X |\psi\rangle$$

$$P_0 = |B|^2$$

$$P_1 = |A|^2$$

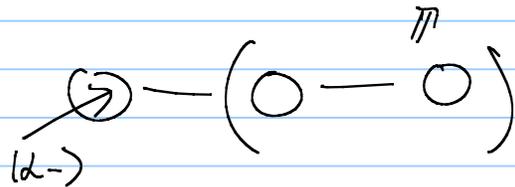
mental  
readjustment

What if on the first measurement we got

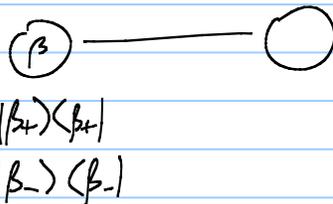
the  $|\alpha-\rangle$  outcome?

$$\langle \alpha- | L_3 \rangle = -i \sin \alpha/2 |0+\rangle + \cos \alpha/2 |1-\rangle \leftarrow \text{BAD}$$

But we wanted:  $-i \sin \alpha/2 |1-\rangle + \cos \alpha/2 |0+\rangle \leftarrow \text{GOOD}$   
 $= \langle \alpha+ | L_3 \rangle$



$$|\psi'\rangle = X_2 Z_3 |\psi_{\text{good}}\rangle$$



Remember  $|\beta_{\pm}\rangle = e^{i\beta/2} |0\rangle \pm e^{-i\beta/2} |1\rangle$

$|\beta_{+}\rangle\langle\beta_{+}| \xleftarrow{X_2} |\psi\rangle$

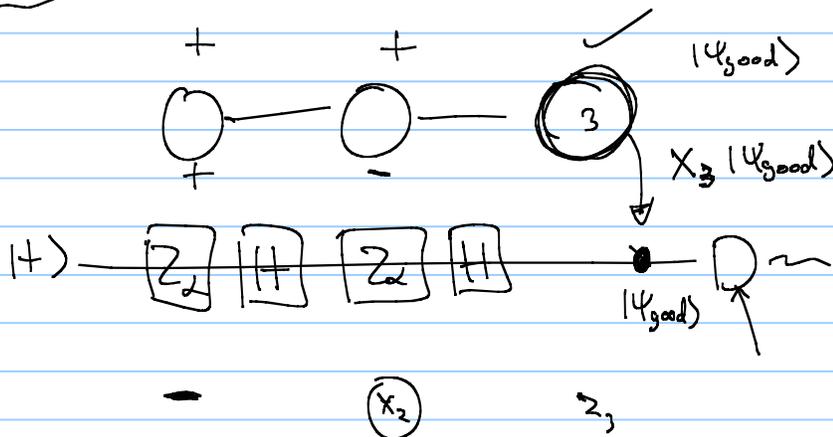
$|\beta_{-}\rangle\langle\beta_{-}|$

Do the measurement

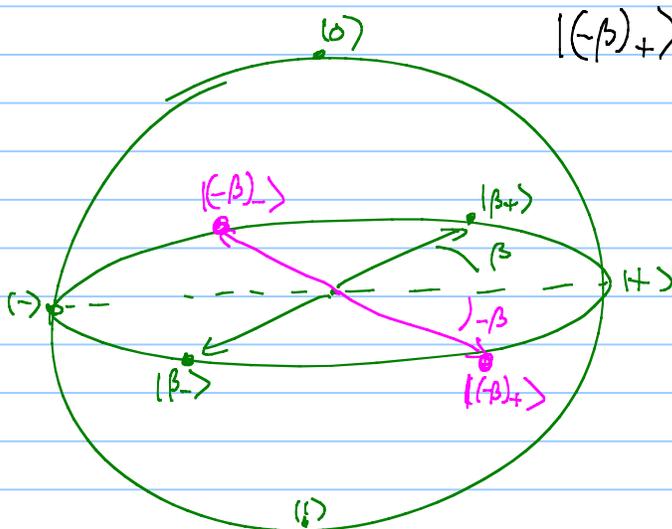
$X|\beta_{+}\rangle\langle\beta_{+}|X$   
 $X|\beta_{-}\rangle\langle\beta_{-}|X$

instead.

$U(A)XAU^{\dagger}$



$$\begin{aligned}
 X|\beta_{+}\rangle &= X \left( e^{i\beta/2} |0\rangle + e^{-i\beta/2} |1\rangle \right) \\
 &= \left( e^{i\beta/2} |1\rangle + e^{-i\beta/2} |0\rangle \right) \\
 &= e^{-i\beta/2} |0\rangle + e^{i\beta/2} |1\rangle
 \end{aligned}$$



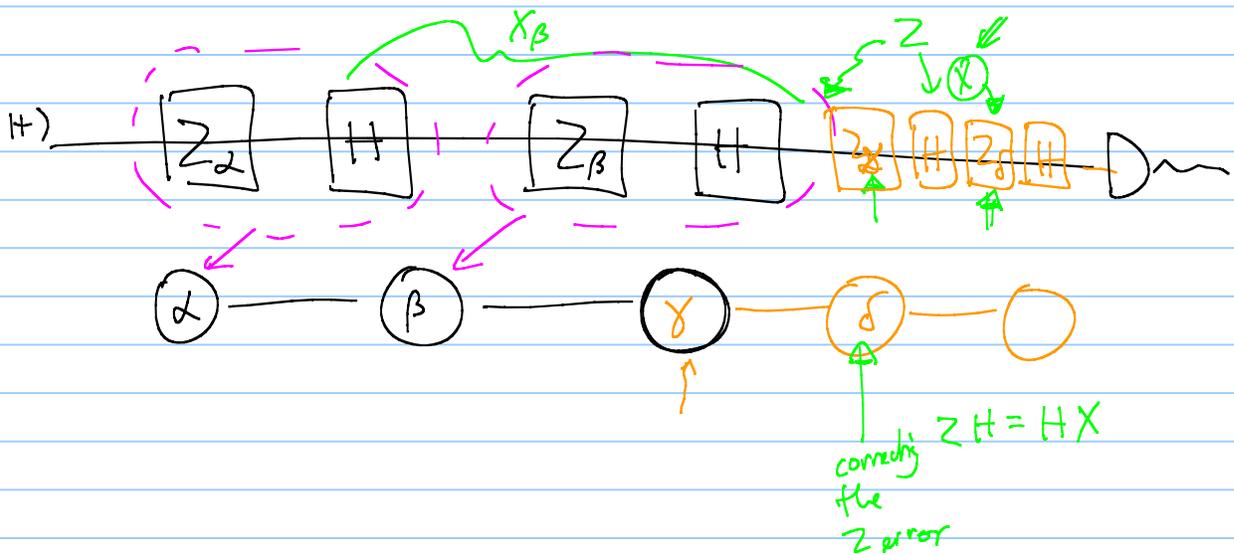
# This is the key idea

Classically feed-forward & change the angle of the second measurement based on the outcome of the first.

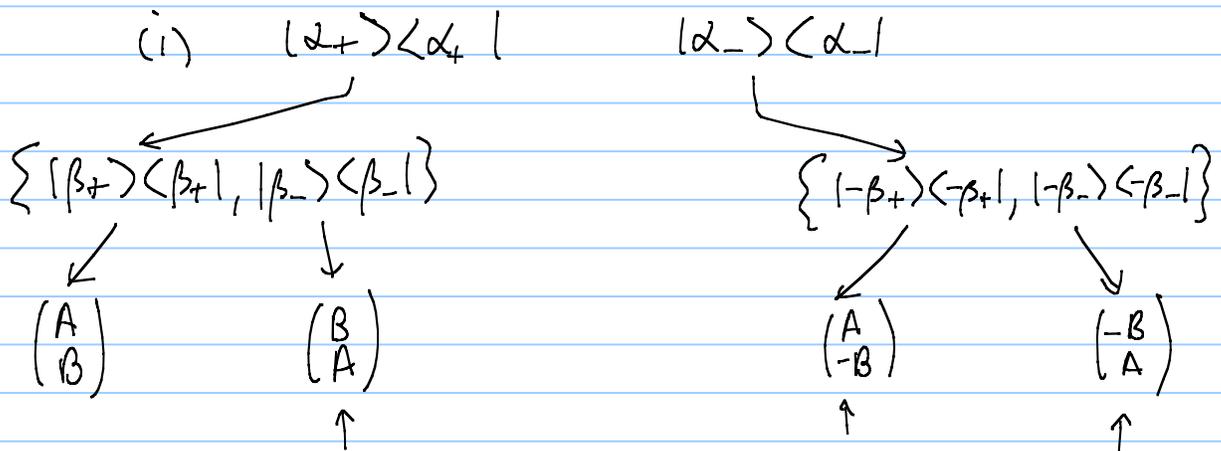
Still have to correct the  $Z_2$  error

$$|\psi''\rangle = \begin{pmatrix} A \\ -B \end{pmatrix}$$

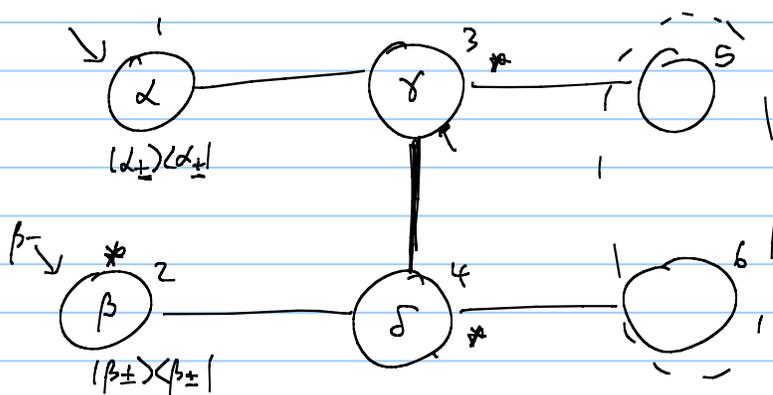
$Z$  error doesn't affect the probabilities  $P_0, P_1$



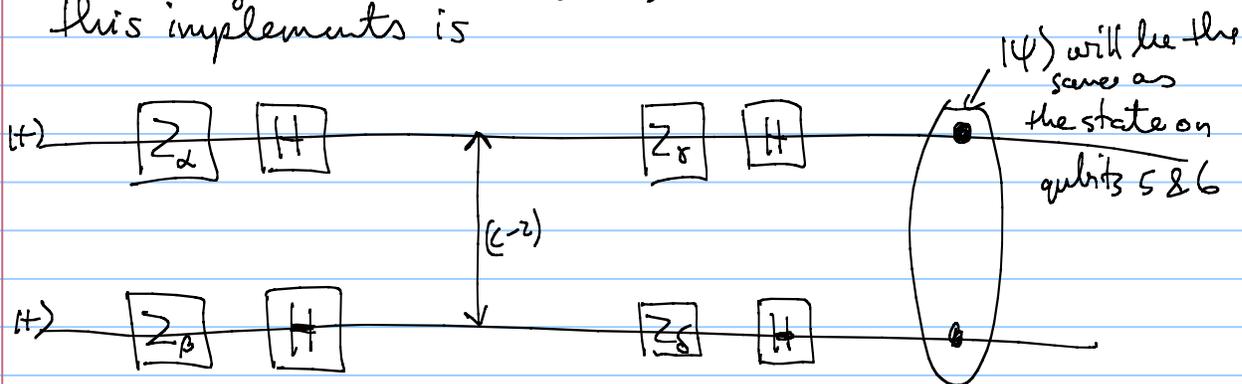
## Summary



## 2-qubit gates?



If we get all "+" (good) outcomes the circuit this implements is

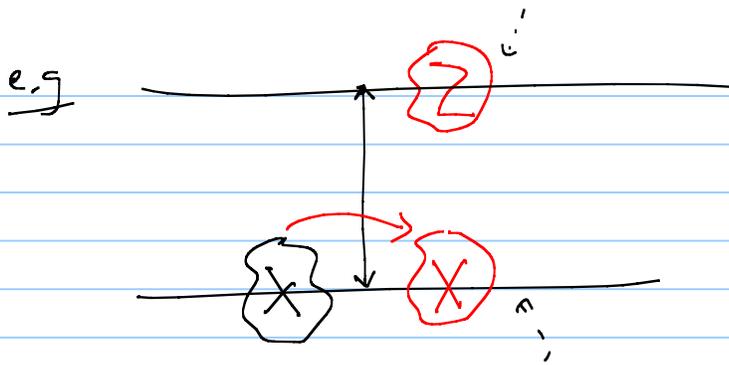


Exercise 1 Prove these outputs are the same.

Exercise 2 Work out what measurements

must be performed on qubits 3 & 4 if we get the  $|\alpha_+\rangle$  outcome on qubit 1 but the  $|\beta_-\rangle$  outcome on qubit 2

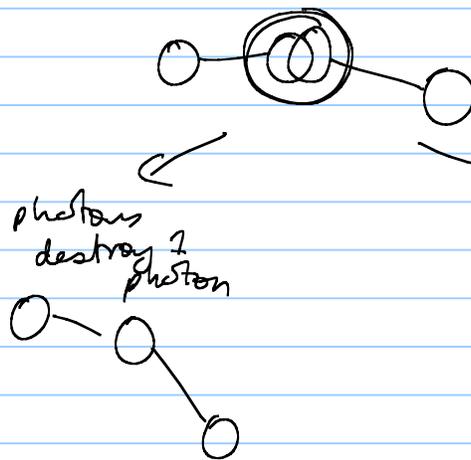
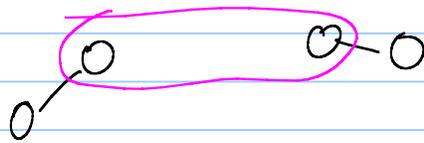
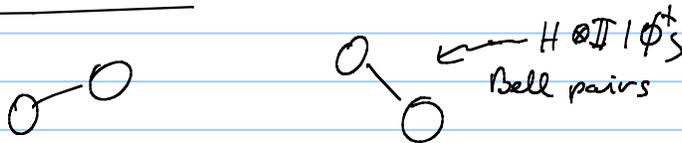
(Hint: Pauli errors going through a CZ gate can "fan out")



$$CZ \otimes X \otimes I = X \otimes Z \otimes CZ$$

## OVERVIEW OF TECHNIQUES FOR BUILDING CLUSTER STATES

### 1/ Fusion



"Redundant Encoding"

$$|0\rangle \Leftrightarrow |00\rangle \Leftrightarrow |000\dots 0\rangle$$

$$|1\rangle \Leftrightarrow |11\rangle \Leftrightarrow |111\dots 1\rangle$$



$$|0\rangle|+\rangle + |1\rangle|-\rangle$$



$$|00\rangle|+\rangle + |11\rangle|-\rangle$$

more generally

$$\frac{(|A\rangle|0\rangle + |A^\perp\rangle|1\rangle)}{\sqrt{2}}$$

nest of the cluster

Bi-orthogonal  
or Schmidt  
decomposition

Redundant Encoding:  $\frac{(|A\rangle|000\rangle + |A^\perp\rangle|111\rangle)}{\sqrt{2}}$

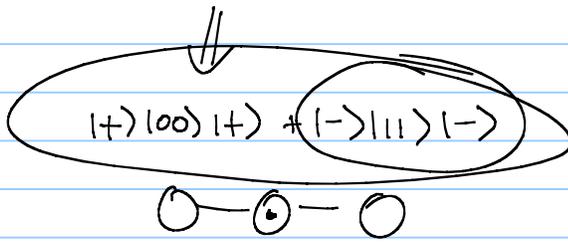
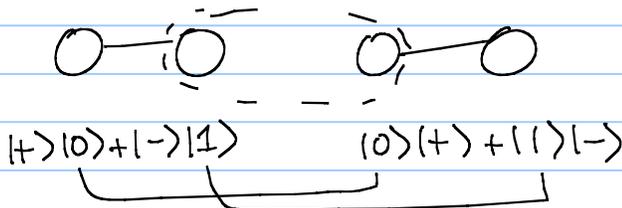
measure in  $|H\rangle$  basis, removes qubits

$$|d_\pm\rangle = e^{i\pi/4}|0\rangle \pm e^{-i\pi/4}|1\rangle$$

$$d \frac{e^{i\pi/4}|000\rangle \pm e^{-i\pi/4}|111\rangle}{\sqrt{2}}$$

not single  
qubit  
measurements

Abstractly



gate fails

$$|V, H\rangle \Rightarrow |2_H\rangle + |2_V\rangle$$

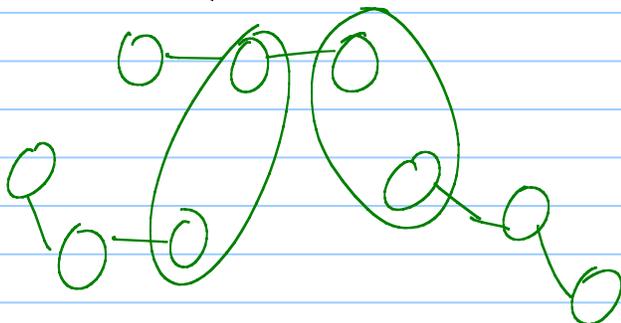
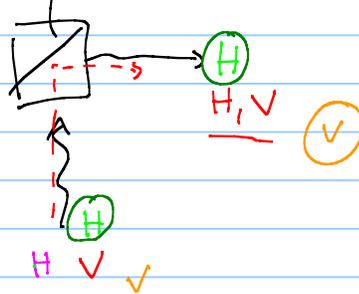
$$\begin{aligned} |H\rangle &\Rightarrow |H\rangle + |V\rangle \\ |V\rangle &\Rightarrow |H\rangle - |V\rangle \end{aligned}$$



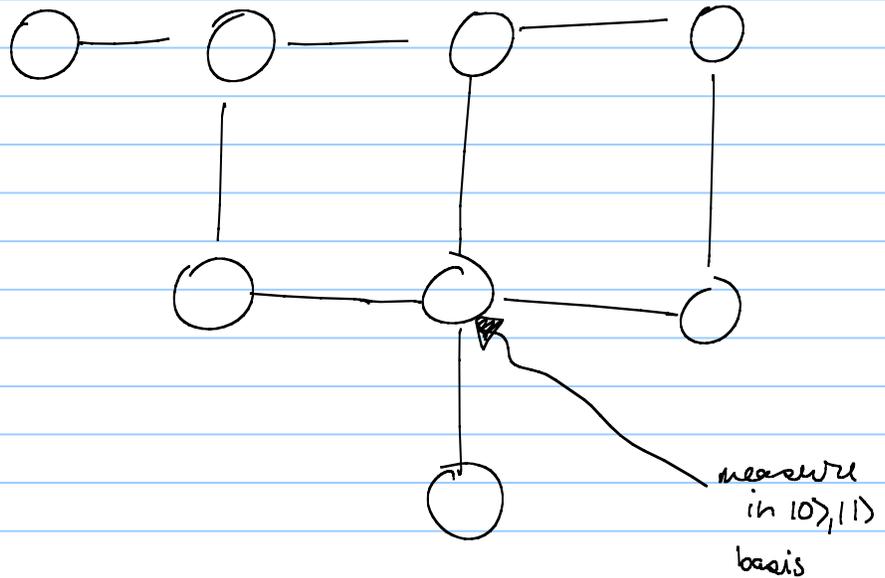
45°

→

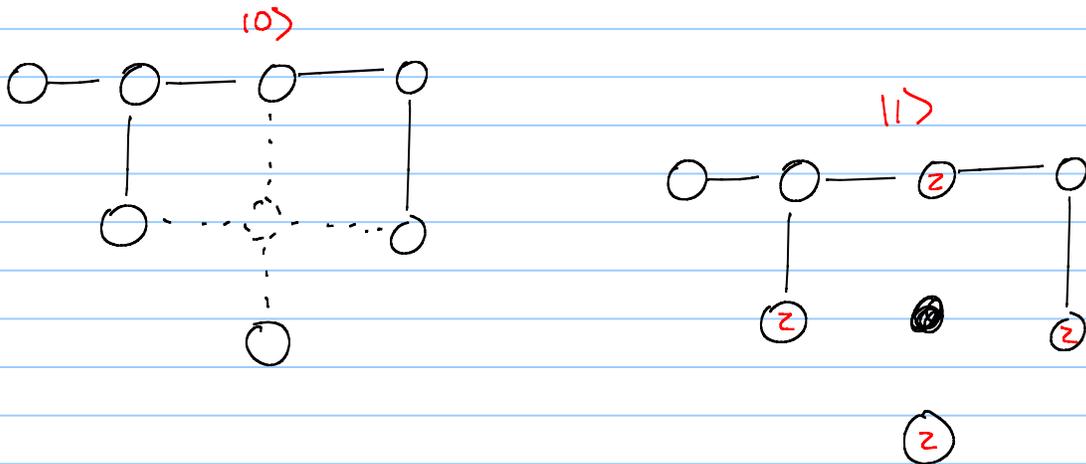
$$\begin{aligned} |H\rangle &= |0\rangle \\ |V\rangle &= |1\rangle \end{aligned}$$



### Exercise 3



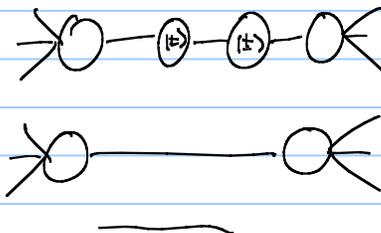
Show that this removes the qubit from the cluster state

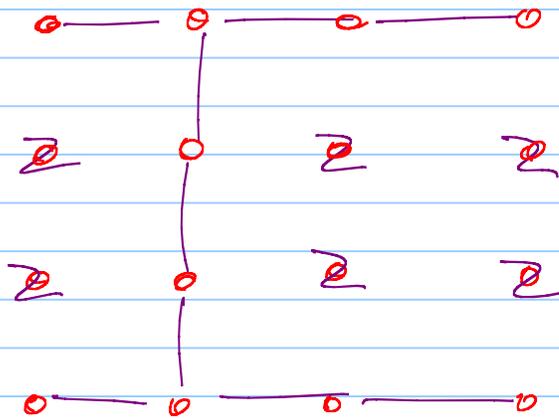


### Exercise 4

Show that two  $X$  basis (i.e.  $|±\rangle$  basis, corresponding to  $\alpha=0$  in  $|\alpha_{±}\rangle$ ) on a linear cluster has the following effect

becomes:





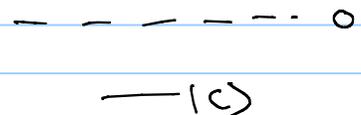
$$0.9 |I\rangle + 0.1 |A\rangle \langle A|$$

## Technique 2

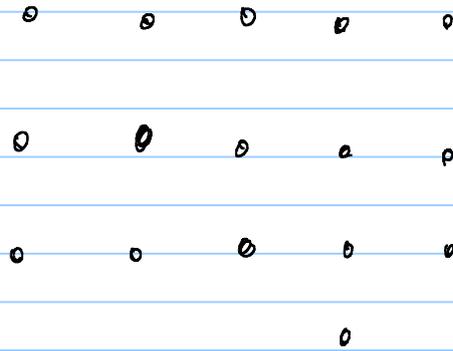
What if nature just gives us cluster states as the ground state of some molecule/crystal/nanotube etc for free?

What would a generic Hamiltonian giving this be?

$$(i) H = -|C\rangle\langle C|$$



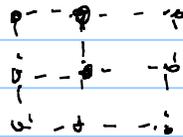
(ii) ? Local Hamiltonian



$\sum$  (products of Pauli operators which are only nontrivial for adjacent qubits)

$$H_{\text{nonlocal}} = X \otimes X \otimes X \otimes \dots \otimes X$$

$H_{\text{local}}$

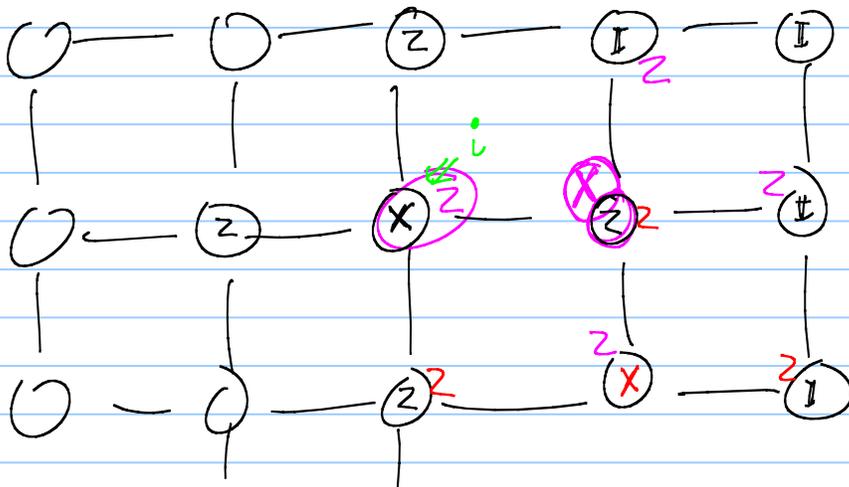


2-body Hamiltonian

$$\sum_{\substack{i,j \\ \uparrow \\ \text{connected} \\ \text{pairs}}} (X_i \otimes X_j + Y_i \otimes Y_j + Z_i \otimes Z_j) \otimes I \otimes I \dots$$

nontrivial Paulis

Consider operators on 5 qubits



local  $\otimes I$   
 $Z \otimes Z \otimes X_i \otimes Z \otimes Z$   
 $K_i$

The following Hamiltonian has the cluster state as its ground state

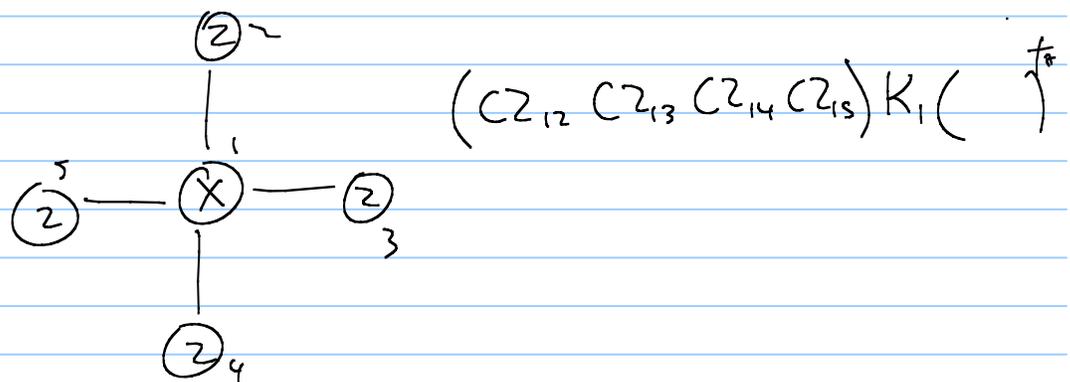
$$H = - \sum_i K_i$$

How to show this? Exercise 5

Think of the cluster as built from the CZ gates

$$|C\rangle = \prod_{\langle ij \rangle} CZ_{ij} |+\dots+\rangle$$

What do the CZ gates do to the  ~~$K_i$~~   $K_i$  operator



Hints  $K_i$  all commute.

Hint The spectrum of the  $K_i$  are all  $\pm 1$