

# Information gain and approximate reversibility of quantum measurements

[see [arXive:quant-ph/0702166v3](https://arxiv.org/abs/quant-ph/0702166v3)]

---

Francesco Buscemi, ERATO-SORST QCI Project, JST



in collaboration with Masahito Hayashi and Michał Horodecki

**IICQI 2007, Kish Island, Iran, September 9th, 2007**

# Overview

- \* Definition of information gain
- \* Definition of disturbance
- \* Balance of information
- \* Tradeoff for general measurements
- \* (Single-outcome analysis)
- \* (Relation with previous proposals:  
Grönwold-Lindblad-Ozawa, Maccone)



# The setting

Let us given an input state  $\rho^Q$   
defined on the input (finite dimensional) Hilbert space  $\mathcal{H}^Q$

# The setting

Let us given an input state  $\rho^Q$   
defined on the input (finite dimensional) Hilbert space  $\mathcal{H}^Q$

Let  $|\Psi^{RQ}\rangle$  be a purification of  $\rho^Q$   
where  $\mathcal{H}^R$  is an auxiliary "reference" system

# The setting

Let us given an input state  $\rho^Q$   
defined on the input (finite dimensional) Hilbert space  $\mathcal{H}^Q$

Let  $|\Psi^{RQ}\rangle$  be a purification of  $\rho^Q$   
where  $\mathcal{H}^R$  is an auxiliary "reference" system

Let the measurement on  $\mathcal{H}^Q$   
be described by the POVM  $\mathbf{P}^Q := \{P_m^Q\}_{m \in \mathcal{X}}$

# Information gain

The measurement on  $Q$  determines an **ensemble decomposition on  $R$** : with probability  $p(m) := \text{Tr}[\rho^Q P_m^Q]$  we observe on  $R$  the conditional state  $\rho_m^{R'} := \text{Tr}_Q[\Psi^{RQ} (\mathbb{1}^R \otimes P_m^Q)]$ .

Our definition of information gain then equals the **Holevo quantity** of the induced ensemble on  $R$ .

# Information gain

The measurement on  $Q$  determines an **ensemble decomposition on  $R$** : with probability  $p(m) := \text{Tr}[\rho^Q P_m^Q]$  we observe on  $R$  the conditional state  $\rho_m^{R'} := \text{Tr}_Q[\Psi^{RQ} (\mathbb{1}^R \otimes P_m^Q)]$ .

Our definition of information gain then equals the **Holevo quantity** of the induced ensemble on  $R$ .

The **(quantum) information gain** is defined to be

$$\iota(\rho^Q, \mathbf{P}^Q) := \sum_m p(m) D(\rho_m^{R'} \parallel \rho^R)$$

# Information gain

The measurement on  $Q$  determines an **ensemble decomposition on  $R$** : with probability  $p(m) := \text{Tr}[\rho^Q P_m^Q]$  we observe on  $R$  the conditional state  $\rho_m^{R'} := \text{Tr}_Q[\Psi^{RQ} (\mathbb{1}^R \otimes P_m^Q)]$ .

Our definition of information gain then equals the **Holevo quantity** of the induced ensemble on  $R$ .

The **(quantum) information gain** is defined to be

$$\begin{aligned} \iota(\rho^Q, \mathbf{P}^Q) &:= \sum_m p(m) D(\rho_m^{R'} \parallel \rho^R) \\ &= S(\rho^R) - \sum_m p(m) S(\rho_m^{R'}) \end{aligned}$$

# Why such a choice?

It is a natural choice for many reasons. In particular:

- it depends **only** on the input state and the POVM
- it is **by construction** positive definite
- it is a natural upper bound to the **classical information gain**, defined as the (classical) mutual information  $I(X : \mathcal{X})$  between the measurement outcomes  $\mathcal{X}$  and the alphabet  $X$  which is eventually encoded in the input state as

$$\rho^Q = \sum_{x \in X} \rho_x^Q$$

# Why such a choice?

It is a natural choice for many reasons. In particular:

- it depends **only** on the input state and the POVM
- it is **by construction** positive definite
- it is a natural upper bound to the **classical information gain**, defined as the (classical) mutual information  $I(X : \mathcal{X})$  between the measurement outcomes  $\mathcal{X}$  and the alphabet  $X$  which is eventually encoded in the input state as

$$\rho^Q = \sum_{x \in X} \rho_x^Q$$

...exactly as Holevo quantity is considered a natural upper bound to the accessible information.

# How to treat disturbance

# How to treat disturbance

In order to analyze the disturbance, the description of the measurement by means of the POVM only is no more sufficient. We have to introduce a **state reduction recipe**, which takes into account the whole statistical description of a quantum measurement, that is, its outcome probability distribution (POVM) as well as its dynamics (state reduction).

# How to treat disturbance

In order to analyze the disturbance, the description of the measurement by means of the POVM only is no more sufficient. We have to introduce a **state reduction recipe**, which takes into account the whole statistical description of a quantum measurement, that is, its outcome probability distribution (POVM) as well as its dynamics (state reduction).

We are led to the notion of **quantum instrument**.

# Quantum instruments

The formalism of **quantum instruments** is the most general setting to describe the full statistics of a quantum measurement. An instrument  $\mathcal{I}^Q$  is defined as follows:

# Quantum instruments

The formalism of **quantum instruments** is the most general setting to describe the full statistics of a quantum measurement. An instrument  $\mathcal{I}^Q$  is defined as follows:

- \* a set of maps  $\{\mathcal{E}_m\}_{m \in \mathcal{X}}$  in one-to-one correspondence with the measurement outcomes is given

# Quantum instruments

The formalism of **quantum instruments** is the most general setting to describe the full statistics of a quantum measurement. An instrument  $\mathcal{I}^Q$  is defined as follows:

- \* a set of maps  $\{\mathcal{E}_m\}_{m \in \mathcal{X}}$  in one-to-one correspondence with the measurement outcomes is given
- \* the probability of obtaining the  $m$ -th outcome is  $p(m) := \text{Tr}[\mathcal{E}_m(\rho^Q)]$

# Quantum instruments

The formalism of **quantum instruments** is the most general setting to describe the full statistics of a quantum measurement. An instrument  $\mathcal{I}^Q$  is defined as follows:

- \* a set of maps  $\{\mathcal{E}_m\}_{m \in \mathcal{X}}$  in one-to-one correspondence with the measurement outcomes is given
- \* the probability of obtaining the  $m$ -th outcome is  $p(m) := \text{Tr}[\mathcal{E}_m(\rho^Q)]$
- \* the “a posteriori” state, given the  $m$ -th outcome, is  $\rho_m^{Q'} := \mathcal{E}_m(\rho^Q)/p(m)$

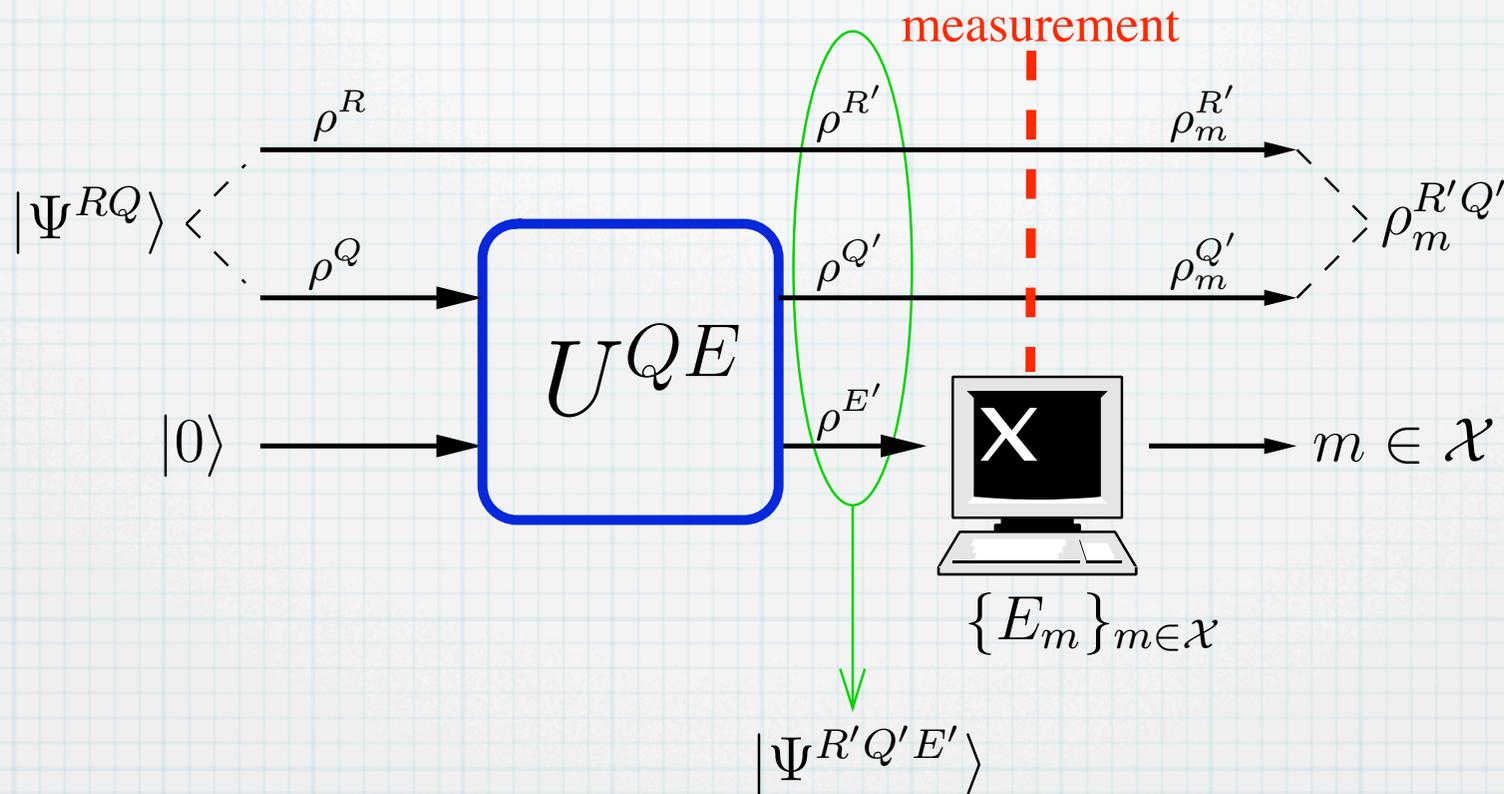
# Indirect measurements

# Indirect measurements

A quantum instrument can be better understood as an **indirect measurement process**

# Indirect measurements

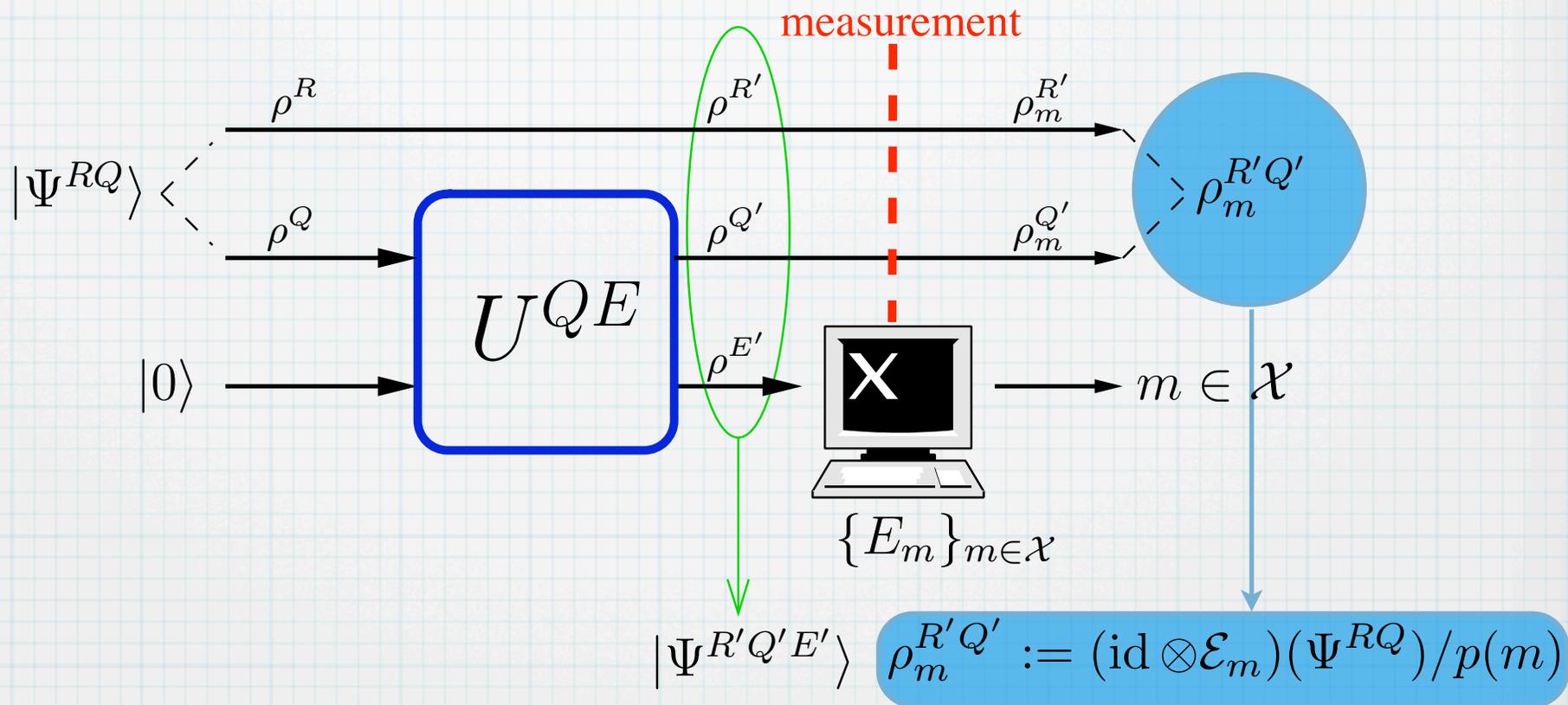
A quantum instrument can be better understood as an **indirect measurement process**



where  $\mathcal{H}^E$  is an **additional ancillary system** and  $\{E_m\}_{m \in \mathcal{X}}$  can be taken as a **Projection-Valued Measure (PVM)**.

# Indirect measurements

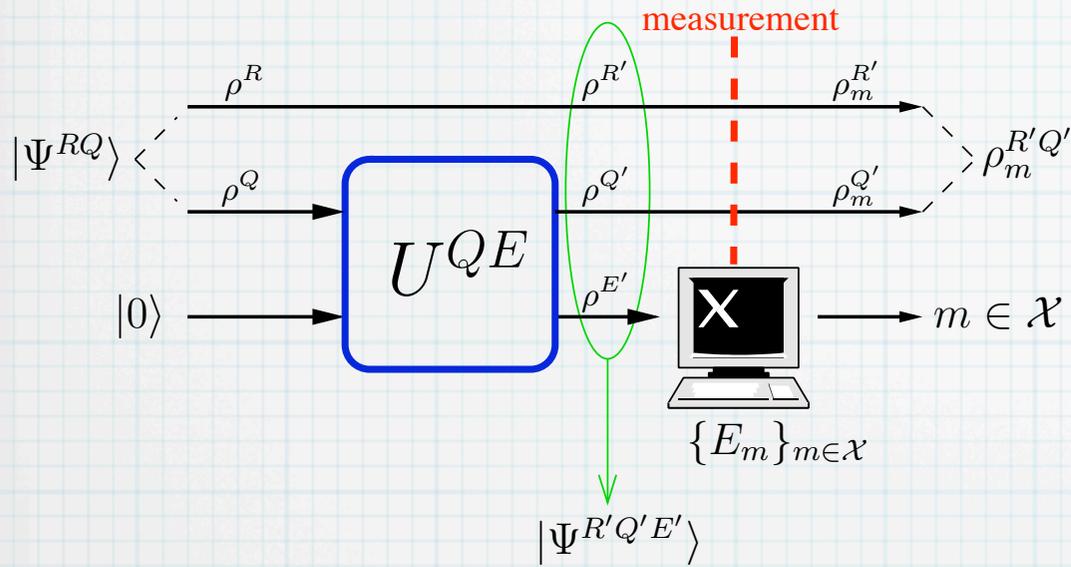
A quantum instrument can be better understood as an **indirect measurement process**



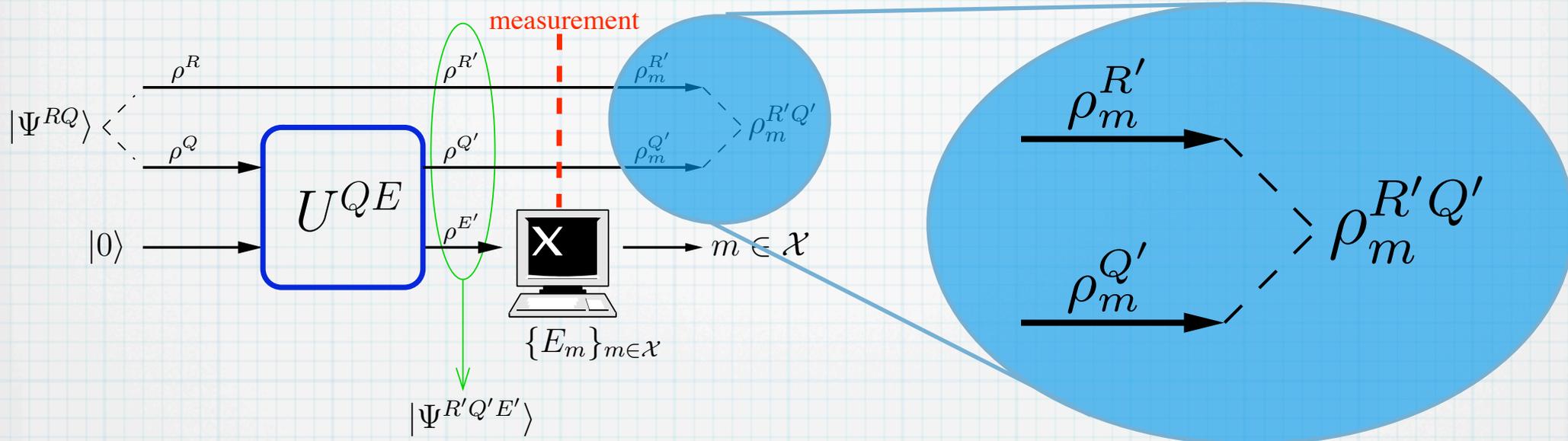
where  $\mathcal{H}^E$  is an **additional ancillary system** and  $\{E_m\}_{m \in \mathcal{X}}$  can be taken as a **Projection-Valued Measure (PVM)**.

# Disturbance

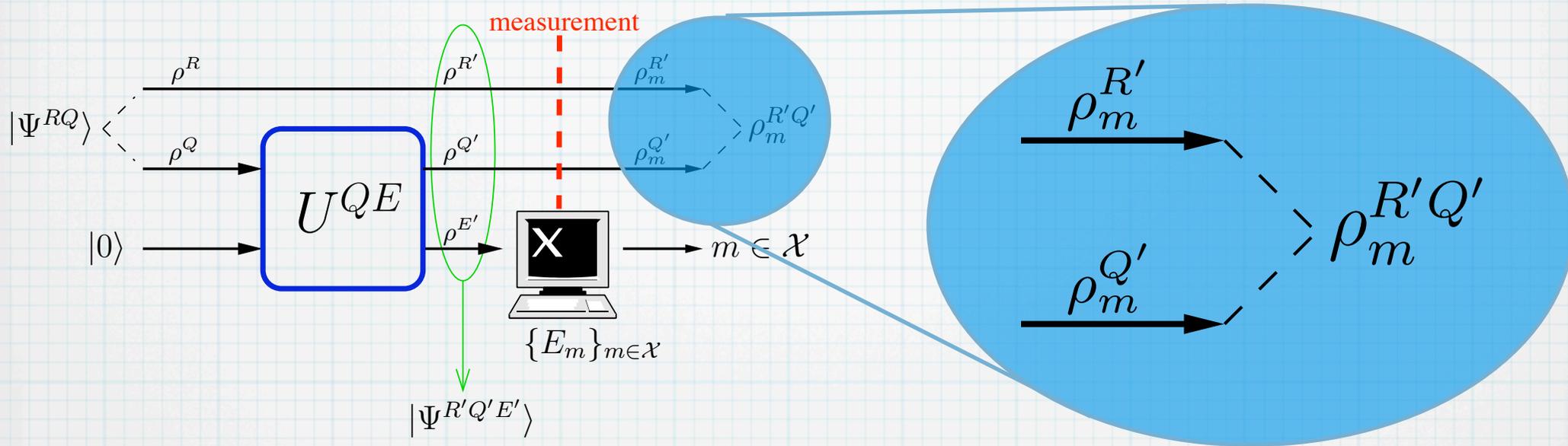
# Disturbance



# Disturbance



# Disturbance



Given the input state and the instrument, we define the **disturbance** as the **conditional coherent information loss**

$$\delta(\rho^Q, \mathcal{I}^Q) := S(\rho^Q) - \sum_m p(m) I_c^{R' \rightarrow Q'}(\rho_m^{R'Q'})$$

where  $I_c^{A \rightarrow B}(\sigma^{AB}) := S(\sigma^B) - S(\sigma^{AB})$

Why such a choice?

# Why such a choice?

Because it is widely accepted that **coherent information** quantifies how well a channel preserves coherence. So, in a measurement process it is natural to consider the same quantity, **conditioned on the outcomes**.

# Why such a choice?

Because it is widely accepted that **coherent information** quantifies how well a channel preserves coherence. So, in a measurement process it is natural to consider the same quantity, **conditioned on the outcomes**.

**Theorem** (generalization of [Schumacher and Westmoreland, QIC (2002)]  
and [Barnum, Nielsen, and Schumacher, PRA (1998)])

# Why such a choice?

Because it is widely accepted that **coherent information** quantifies how well a channel preserves coherence. So, in a measurement process it is natural to consider the same quantity, **conditioned on the outcomes**.

**Theorem** (generalization of [Schumacher and Westmoreland, QIC (2002)] and [Barnum, Nielsen, and Schumacher, PRA (1998)])

there exist channels  $\{\mathcal{R}_m\}_{m \in \mathcal{X}}$  such that

$$F^e(\rho^Q, \sum_m \mathcal{R}_m \circ \mathcal{E}_m) \geq 1 - \sqrt{2\delta(\rho^Q, \mathcal{I}^Q)}$$

# Why such a choice?

Because it is widely accepted that **coherent information** quantifies how well a channel preserves coherence. So, in a measurement process it is natural to consider the same quantity, **conditioned on the outcomes**.

**Theorem** (generalization of [Schumacher and Westmoreland, QIC (2002)] and [Barnum, Nielsen, and Schumacher, PRA (1998)])

☑ there exist channels  $\{\mathcal{R}_m\}_{m \in \mathcal{X}}$  such that

$$F^e(\rho^Q, \sum_m \mathcal{R}_m \circ \mathcal{E}_m) \geq 1 - \sqrt{2\delta(\rho^Q, \mathcal{I}^Q)}$$

☑ given a set of channels  $\{\mathcal{R}_m\}_{m \in \mathcal{X}}$  it holds that

$$\delta(\rho^Q, \mathcal{I}^Q) \leq h \left( 1 - F^e(\rho^Q, \sum_m \mathcal{R}_m \circ \mathcal{E}_m) \right)$$

# Output of a measurement

# Output of a measurement

The global **reference+system+apparatus** state after the measurement can be written w.l.o.g. as follows

$$\Upsilon^{R'Q'E'} \mathcal{X} := \sum_m p(m) \Psi_m^{R'Q'E'} \otimes |m\rangle\langle m| \mathcal{X}$$

# Output of a measurement

The global **reference+system+apparatus** state after the measurement can be written w.l.o.g. as follows

$$\Upsilon^{R'Q'E'} \mathcal{X} := \sum_m p(m) \Psi_m^{R'Q'E'} \otimes |m\rangle \langle m| \mathcal{X}$$

where  $|\Psi_m^{R'Q'E'}\rangle := (\mathbb{1}^{R'Q'} \otimes E_m^{1/2})(\mathbb{1}^R \otimes U^{QE})(|\Psi^{RQ}\rangle \otimes |0^E\rangle) / \sqrt{p(m)}$

and correspondingly  $\rho_m^{R'Q'} = \text{Tr}_{E'}[\Psi_m^{R'Q'E'}]$

# Output of a measurement

The global **reference+system+apparatus** state after the measurement can be written w.l.o.g. as follows

$$\Upsilon^{R'Q'E'} \mathcal{X} := \sum_m p(m) \Psi_m^{R'Q'E'} \otimes |m\rangle \langle m| \mathcal{X}$$

$E_m$  rank one implies  $E_m^{1/2}$  rank one

where  $|\Psi_m^{R'Q'E'}\rangle := (\mathbb{1}^{R'Q'} \otimes E_m^{1/2}) (\mathbb{1}^R \otimes U^{QE}) (|\Psi^{RQ}\rangle \otimes |0^E\rangle) / \sqrt{p(m)}$

and correspondingly  $\rho_m^{R'Q'} = \text{Tr}_{E'} [\Psi_m^{R'Q'E'}]$

# Output of a measurement

The global **reference+system+apparatus** state after the measurement can be written w.l.o.g. as follows

$$\Upsilon^{R'Q'E'} \mathcal{X} := \sum_m p(m) \Psi_m^{R'Q'E'} \otimes |m\rangle \langle m| \mathcal{X}$$

$E_m$  rank one implies  $E_m^{1/2}$  rank one

where  $|\Psi_m^{R'Q'E'}\rangle := (\mathbb{1}^{R'Q'} \otimes E_m^{1/2}) (\mathbb{1}^R \otimes U^{QE}) (|\Psi^{RQ}\rangle \otimes |0^E\rangle) / \sqrt{p(m)}$

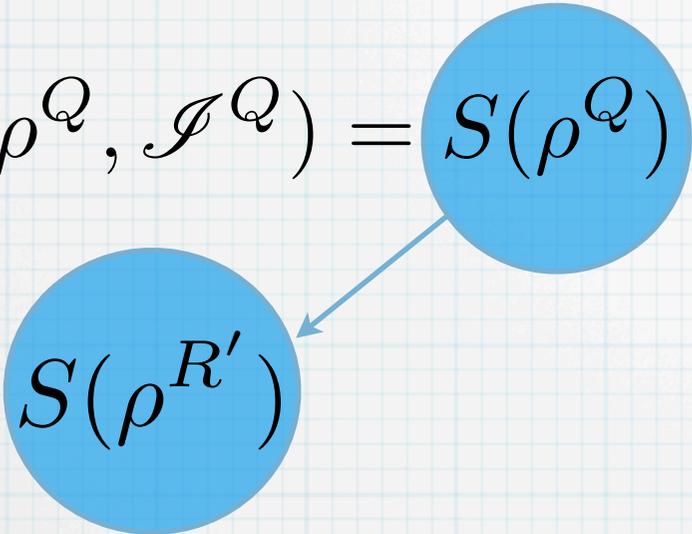
and correspondingly  $\rho_m^{R'Q'} = \text{Tr}_{E'} [\Psi_m^{R'Q'E'}]$

which in turn implies that  $\rho_m^{R'Q'}$  is pure, that means, reference+system, given the  $m$ -th outcome, is factorized from the rest of the universe

# Rewriting disturbance

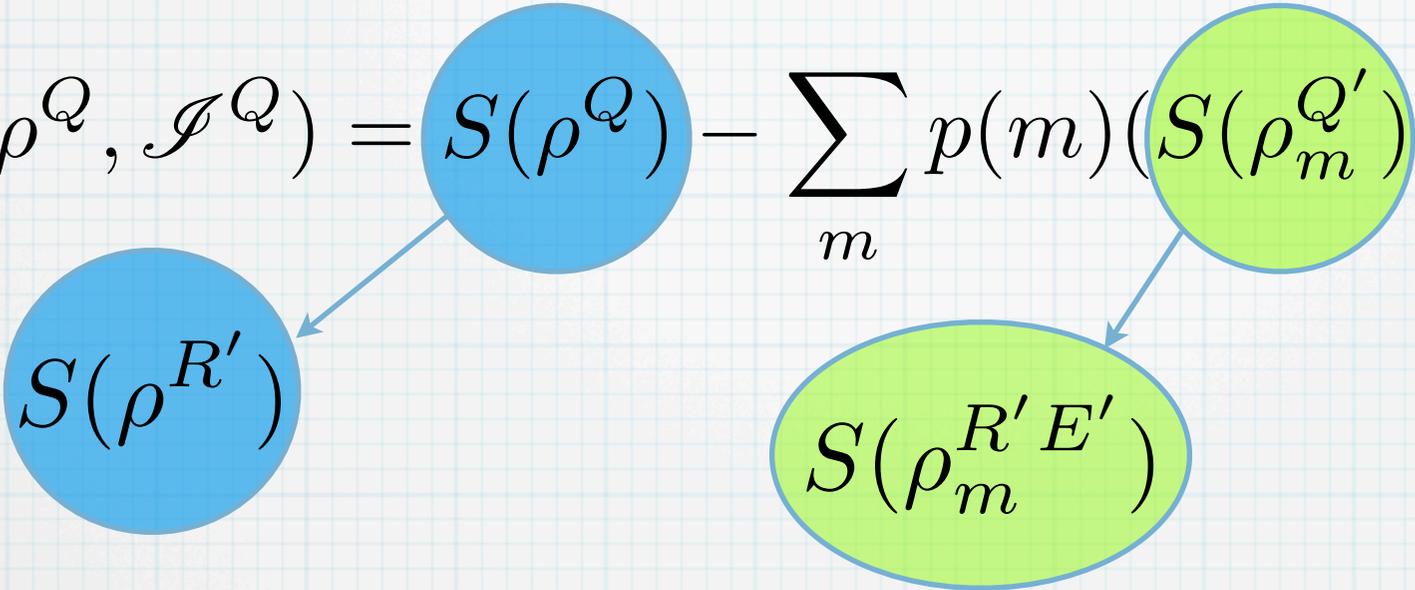
$$\delta(\rho^Q, \mathcal{I}^Q) = S(\rho^Q) - \sum_m p(m) (S(\rho_m^{Q'}) - S(\rho_m^{R'Q'}))$$

# Rewriting disturbance

$$\delta(\rho^Q, \mathcal{I}^Q) = S(\rho^Q) - \sum_m p(m) (S(\rho_m^{Q'}) - S(\rho_m^{R'Q'}))$$


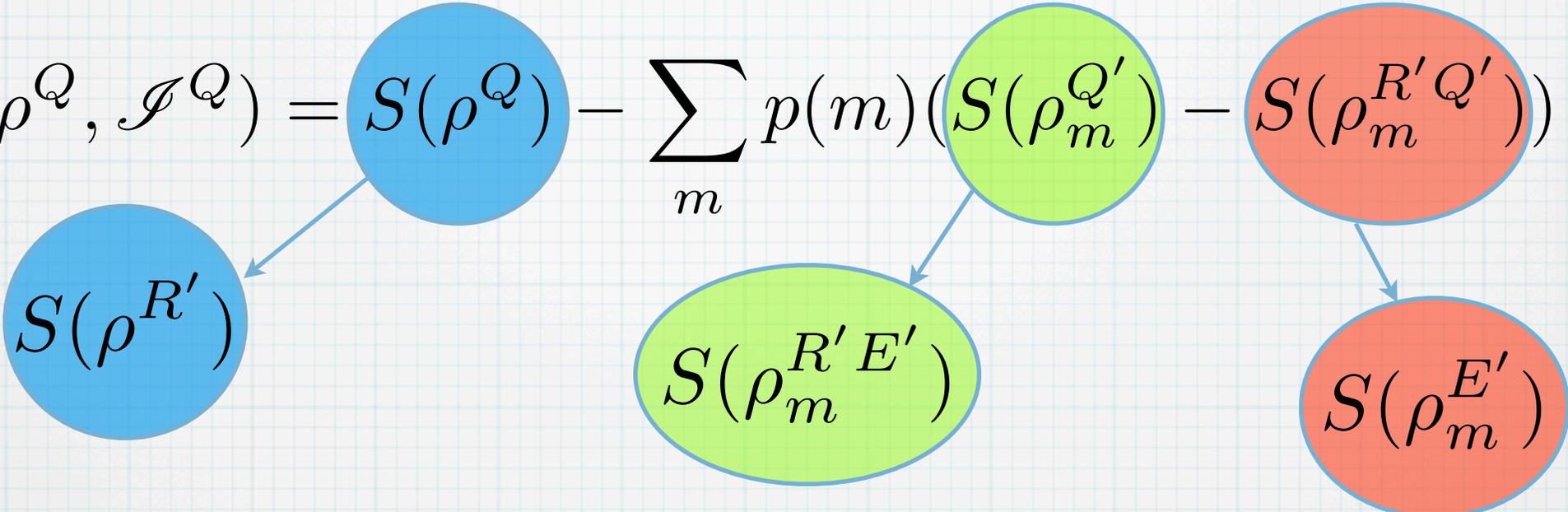
A diagram consisting of two blue circles. The upper circle contains the mathematical expression  $S(\rho^Q)$ . A light blue arrow points from this circle down and to the left towards a second blue circle. The lower circle contains the mathematical expression  $S(\rho^{R'})$ .

# Rewriting disturbance

$$\delta(\rho^Q, \mathcal{I}^Q) = S(\rho^Q) - \sum_m p(m) (S(\rho_m^{Q'}) - S(\rho_m^{R'Q'}))$$


The diagram illustrates the decomposition of the disturbance term. The first term,  $S(\rho^Q)$ , is shown in a blue circle. An arrow points from this circle to another blue circle containing  $S(\rho^{R'})$ . The second term,  $S(\rho_m^{Q'})$ , is shown in a green circle. An arrow points from this circle to a green oval containing  $S(\rho_m^{R'E'})$ .

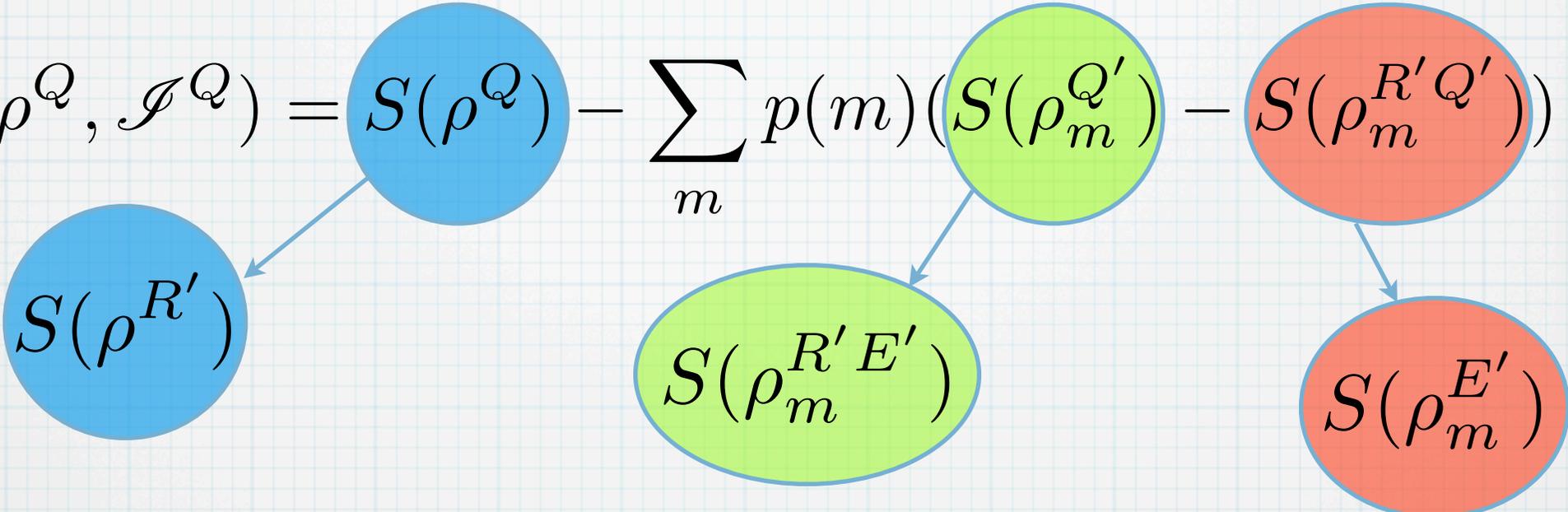
# Rewriting disturbance

$$\delta(\rho^Q, \mathcal{I}^Q) = S(\rho^Q) - \sum_m p(m) (S(\rho_m^{Q'}) - S(\rho_m^{R'Q'}))$$


The diagram illustrates the decomposition of the disturbance term in the equation above. It consists of six nodes and three arrows:

- A blue circle containing  $S(\rho^Q)$  has an arrow pointing to a blue circle containing  $S(\rho^{R'})$ .
- A green circle containing  $S(\rho_m^{Q'})$  has an arrow pointing to a green oval containing  $S(\rho_m^{R'E'})$ .
- A red circle containing  $S(\rho_m^{R'Q'})$  has an arrow pointing to a red circle containing  $S(\rho_m^{E'})$ .

# Rewriting disturbance

$$\delta(\rho^Q, \mathcal{I}^Q) = S(\rho^Q) - \sum_m p(m) (S(\rho_m^{Q'}) - S(\rho_m^{R'Q'}))$$


The diagram illustrates the decomposition of the disturbance equation. The main equation is at the top. Below it, three terms are shown in colored circles: a blue circle for  $S(\rho^{R'})$ , a green circle for  $S(\rho_m^{R'E'})$ , and a red circle for  $S(\rho_m^{E'})$ . Arrows point from the main equation to these three terms: from  $S(\rho^Q)$  to  $S(\rho^{R'})$ , from  $S(\rho_m^{Q'})$  to  $S(\rho_m^{R'E'})$ , and from  $S(\rho_m^{R'Q'})$  to  $S(\rho_m^{E'})$ .

In other words, the disturbance equals the **total correlations between the reference and the apparatus**:

$$\delta(\rho^Q, \mathcal{I}^Q) = I^{R':E'} \chi(\Upsilon^{R'E'} \chi)$$

where  $\Upsilon^{R'E'} \chi := \text{Tr}_{Q'}[\Upsilon^{R'Q'E'} \chi]$

# Tradeoff relation

# Tradeoff relation

By noticing that  $\iota(\rho^Q, \mathbf{P}^Q) = I^{R':\mathcal{X}}(\Upsilon^{R'}\mathcal{X})$

# Tradeoff relation

By noticing that  $I(\rho^Q, \mathbf{P}^Q) = I^{R':\mathcal{X}}(\Upsilon^{R'}\mathcal{X})$

and by the **chain rule** for the quantum mutual information

$$I^{A:B}(\sigma^{AB}) + I^{A:C|B}(\sigma^{ABC}) = I^{A:BC}(\sigma^{ABC})$$

# Tradeoff relation

By noticing that  $\iota(\rho^Q, \mathbf{P}^Q) = I^{R':\mathcal{X}}(\Upsilon^{R'\mathcal{X}})$

and by the **chain rule** for the quantum mutual information

$$I^{A:B}(\sigma^{AB}) + I^{A:C|B}(\sigma^{ABC}) = I^{A:BC}(\sigma^{ABC})$$

we arrive at the following

$$\iota(\rho^Q, \mathbf{P}^Q) + \Delta(\rho^Q, \mathcal{I}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

# Tradeoff relation

By noticing that  $\iota(\rho^Q, \mathbf{P}^Q) = I^{R':\mathcal{X}}(\Upsilon^{R'\mathcal{X}})$

and by the **chain rule** for the quantum mutual information

$$I^{A:B}(\sigma^{AB}) + I^{A:C|B}(\sigma^{ABC}) = I^{A:BC}(\sigma^{ABC})$$

we arrive at the following

$$\iota(\rho^Q, \mathbf{P}^Q) + \Delta(\rho^Q, \mathcal{I}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

Information

# Tradeoff relation

By noticing that  $\iota(\rho^Q, \mathbf{P}^Q) = I^{R':\mathcal{X}}(\Upsilon^{R'\mathcal{X}})$

and by the **chain rule** for the quantum mutual information

$$I^{A:B}(\sigma^{AB}) + I^{A:C|B}(\sigma^{ABC}) = I^{A:BC}(\sigma^{ABC})$$

we arrive at the following

$$\iota(\rho^Q, \mathbf{P}^Q) + \Delta(\rho^Q, \mathcal{I}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

Information

Disturbance

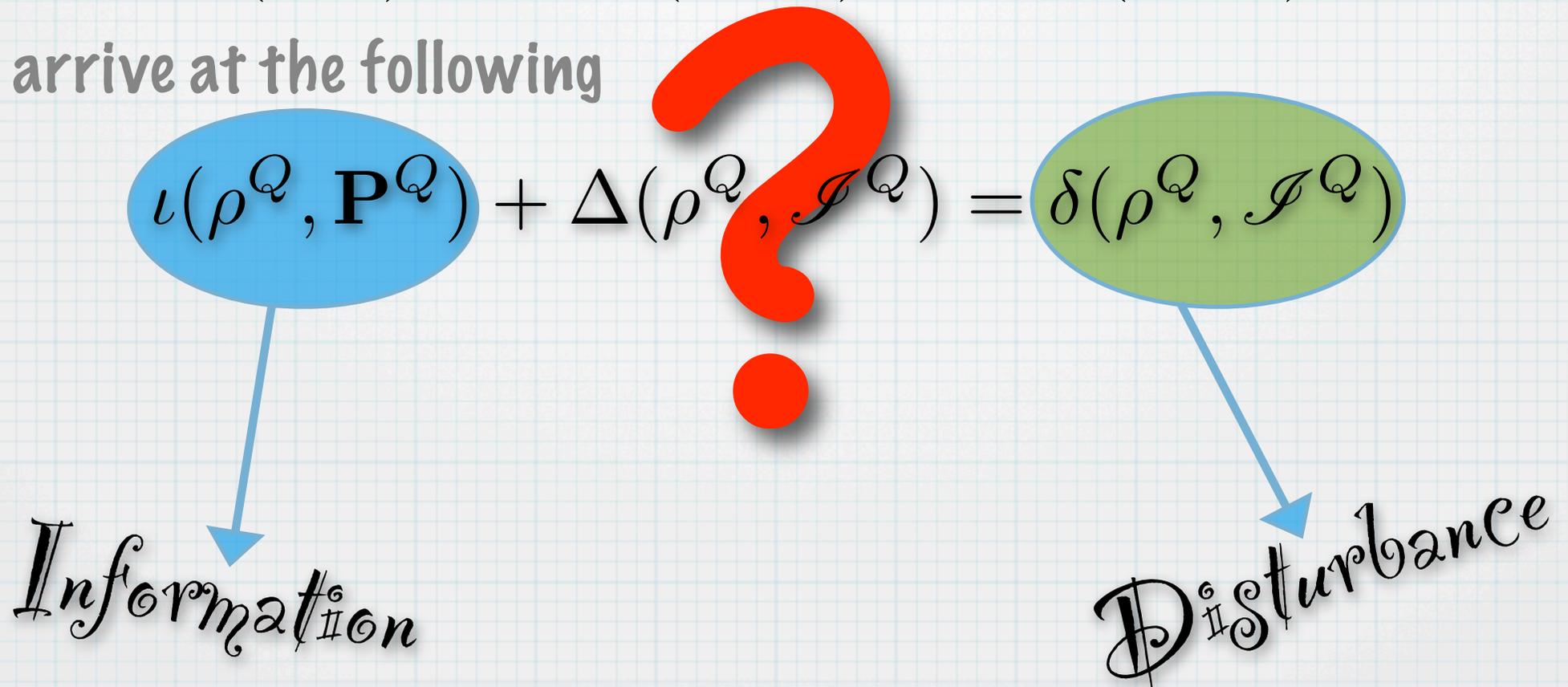
# Tradeoff relation

By noticing that  $\iota(\rho^Q, \mathbf{P}^Q) = I^{R':\mathcal{X}}(\Upsilon^{R'\mathcal{X}})$

and by the **chain rule** for the quantum mutual information

$$I^{A:B}(\sigma^{AB}) + I^{A:C|B}(\sigma^{ABC}) = I^{A:BC}(\sigma^{ABC})$$

we arrive at the following



# Meaning of $\Delta$

$$\Delta(\rho^Q, \mathcal{I}^Q) := I^{R':E'|\mathcal{X}}(\Upsilon^{R'E'\mathcal{X}})$$

# Meaning of $\Delta$

$$\Delta(\rho^Q, \mathcal{I}^Q) := I^{R':E'|\mathcal{X}}(\Upsilon^{R'E'}\mathcal{X})$$

$$= \sum_m p(m) D(\rho_m^{R'E'} \parallel \rho_m^{R'} \otimes \rho_m^{E'})$$

# Meaning of $\Delta$

$$\Delta(\rho^Q, \mathcal{I}^Q) := I^{R':E'|\mathcal{X}}(\Upsilon^{R'E'}\mathcal{X})$$

$$= \sum_m p(m) D(\rho_m^{R'E'} \parallel \rho_m^{R'} \otimes \rho_m^{E'})$$

namely, it represents the **average amount of correlations** between the reference and the internal degrees of freedom of the apparatus.

# Meaning of $\Delta$

$$\begin{aligned}\Delta(\rho^Q, \mathcal{I}^Q) &:= I^{R':E'|\mathcal{X}}(\Upsilon^{R'E'}\mathcal{X}) \\ &= \sum_m p(m) D(\rho_m^{R'E'} \parallel \rho_m^{R'} \otimes \rho_m^{E'})\end{aligned}$$

namely, it represents the **average amount of correlations** between the reference and the internal degrees of freedom of the apparatus.

Usually, apparatus internal degrees of freedom are **out of our control**, and the information gain is strictly less than the disturbance introduced.

# Hidden correlations

# Hidden correlations

Generally, hidden degrees of freedom of the apparatus remain entangled with the system, in such a way that the information extracted is less than the disturbance introduced.

# Hidden correlations

Generally, hidden degrees of freedom of the apparatus remain entangled with the system, in such a way that the information extracted is less than the disturbance introduced.

However, if the instrument is “single-Kraus”, or “multiplicity free”, that is  $\mathcal{E}_m(\rho^Q) = T_m \rho^Q T_m^\dagger$  for all  $m$ , then  $\Delta=0$ , and the tradeoff relation simplifies as follows

# Hidden correlations

Generally, hidden degrees of freedom of the apparatus remain entangled with the system, in such a way that the information extracted is less than the disturbance introduced.

However, if the instrument is “single-Kraus”, or “multiplicity free”, that is  $\mathcal{E}_m(\rho^Q) = T_m \rho^Q T_m^\dagger$  for all  $m$ , then  $\Delta=0$ , and the tradeoff relation simplifies as follows

$$\iota(\rho^Q, \mathbf{P}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

# Hidden correlations

Generally, hidden degrees of freedom of the apparatus remain entangled with the system, in such a way that the information extracted is less than the disturbance introduced.

However, if the instrument is “**single-Kraus**”, or “**multiplicity free**”, that is  $\mathcal{E}_m(\rho^Q) = T_m \rho^Q T_m^\dagger$  for all  $m$ , then  $\Delta=0$ , and the tradeoff relation simplifies as follows

$$\iota(\rho^Q, \mathbf{P}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

Equivalently, this is the case in which the PVM operators  $E_m$  are **all rank-one**.

# Hidden correlations

Generally, **hidden degrees of freedom** of the apparatus remain entangled with the system, in such a way that the information extracted is less than the disturbance introduced.

However, if the instrument is “**single-Kraus**”, or “**multiplicity free**”, that is  $\mathcal{E}_m(\rho^Q) = T_m \rho^Q T_m^\dagger$  for all  $m$ , then  $\Delta=0$ , and the tradeoff relation simplifies as follows

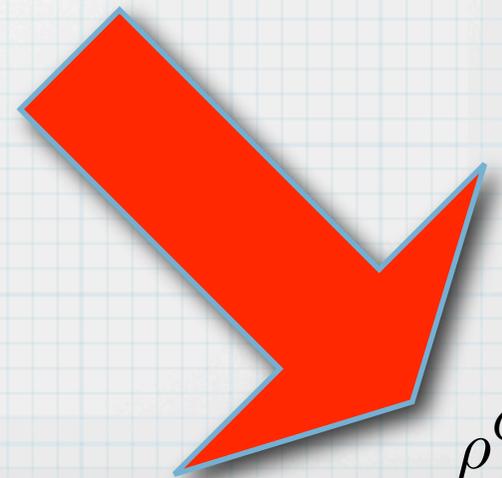
$$\iota(\rho^Q, \mathbf{P}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

Equivalently, this is the case in which the PVM operators  $E_m$  are **all rank-one**.

internal degrees of freedom  
+  
environment interacting with these

Wigner's friends...

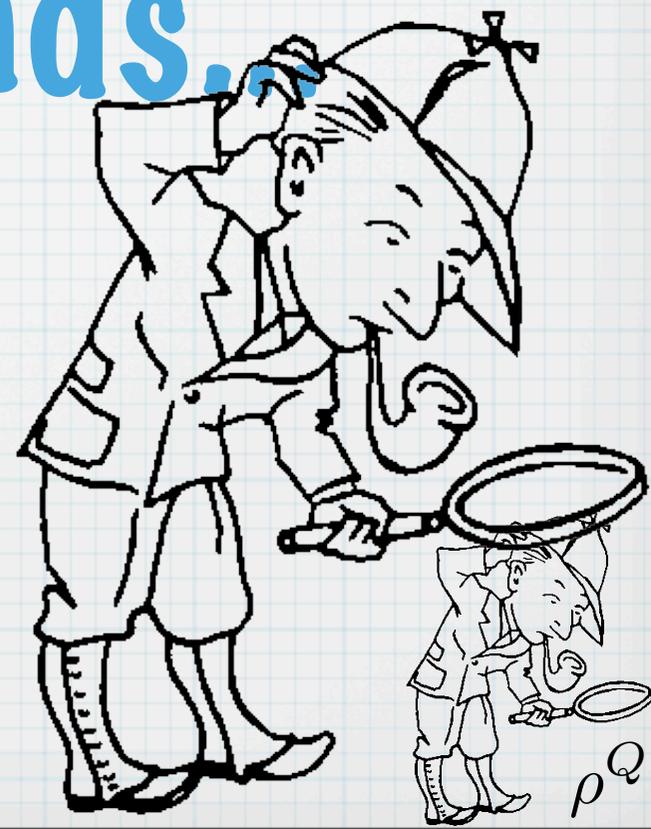
Wigner's friends...



# Wigner's friends...



# Wigner's friends.



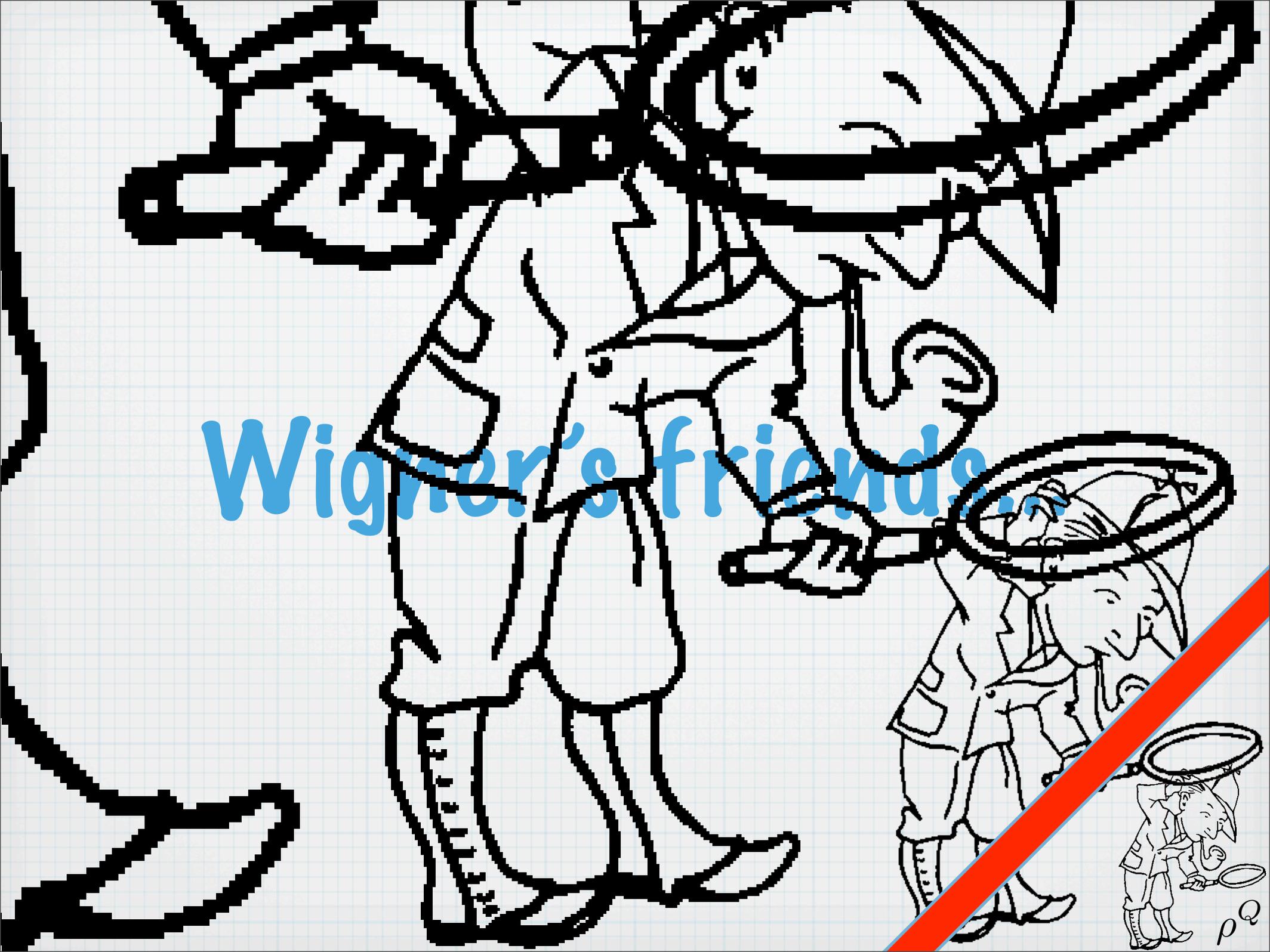
# Wigner's friends



Wigner's friends



# Wigner's friends



Conclusion: general  
tradeoff

# Conclusion: general tradeoff

$$\iota(\rho^Q, \mathbf{P}^Q) + \Delta(\rho^Q, \mathcal{I}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

# Conclusion: general tradeoff

$$I(\rho^Q, \mathbf{P}^Q) + \Delta(\rho^Q, \mathcal{I}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

Information

Disturbance

# Conclusion: general tradeoff

$$I(\rho^Q, \mathbf{P}^Q) + \Delta(\rho^Q, \mathcal{I}^Q) = \delta(\rho^Q, \mathcal{I}^Q)$$

Information

Hidden correlations

Disturbance



kyami!

